# EE540 Advance Electromagnetic Theory & Antennas

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- Far Fields of aperture antenna
- Step 1:
  - Establish the EM fields specified on the aperture surface
- Step 2:
  - Calculate the 2-D FT of the aperture fields
- Step 3:
  - Find the far field electric fields from equations (24.1) and (24.2),
  - Far field magnetic fields can be found from the electric field (if required) from  $\vec{H}_{ff} = \frac{\hat{r} \times \vec{E}_{ff}}{\eta}$

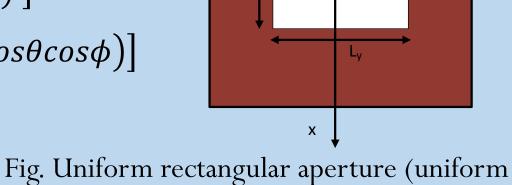
• Radiation from Uniform rectangular aperture

• 
$$\vec{E}_a = E_{ay}\hat{y} = \begin{cases} E_0\hat{y}, & |x| \le \frac{L_x}{2}, |y| \le \frac{L_y}{2}, z = 0\\ 0, & otherwise \end{cases}$$

• Hence, from equations (24.1) and (24.2), we have,

• 
$$E_{ff\theta} \cong 2 \frac{j\beta e^{(-j\beta r)}}{4\pi r} [(\tilde{E}_{ay}sin\phi)]$$
  
and  $E_{ff\phi} \cong 2 \frac{j\beta e^{(-j\beta r)}}{4\pi r} [(\tilde{E}_{ay}cos\theta cos\phi)]$ 

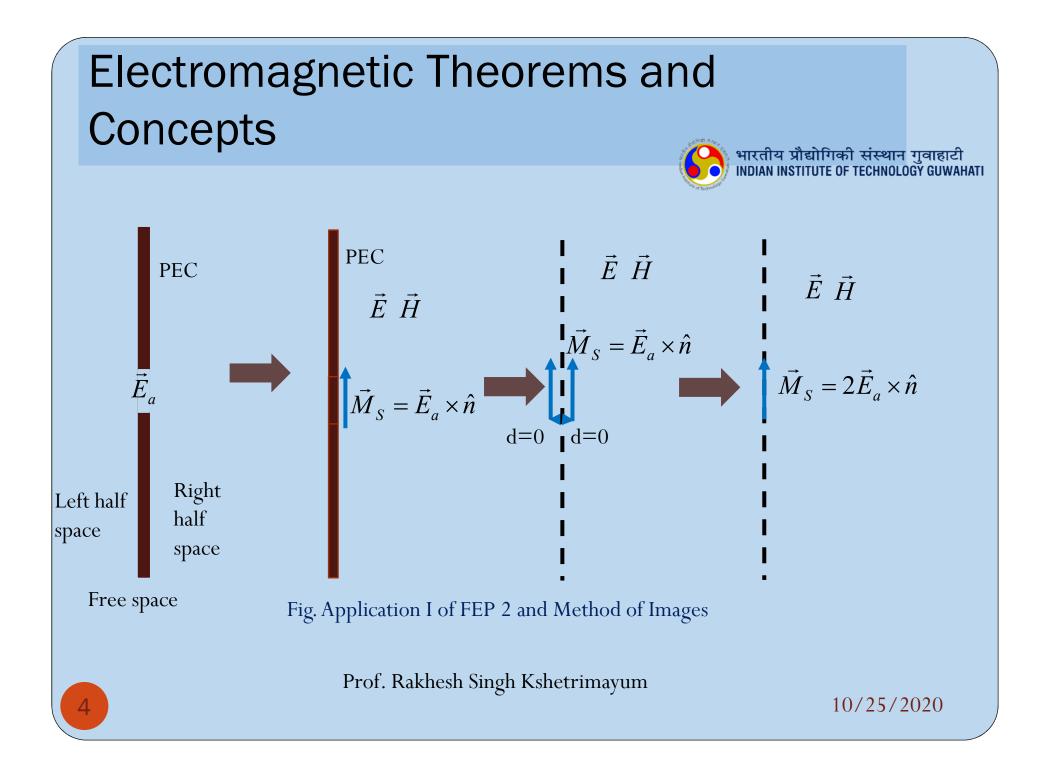
• Why this factor of 2?



7

amplitude and phase of the aperture field)

y



• Let us find the 2 – D FT of the aperture field,

• 
$$\tilde{E}_{ay} = \int_{-\frac{L_y}{2}}^{\frac{L_y}{2}} \int_{-\frac{L_x}{2}}^{\frac{L_x}{2}} E_0 e^{(j\beta_x x' + j\beta_y y')} dx' dy'$$
  

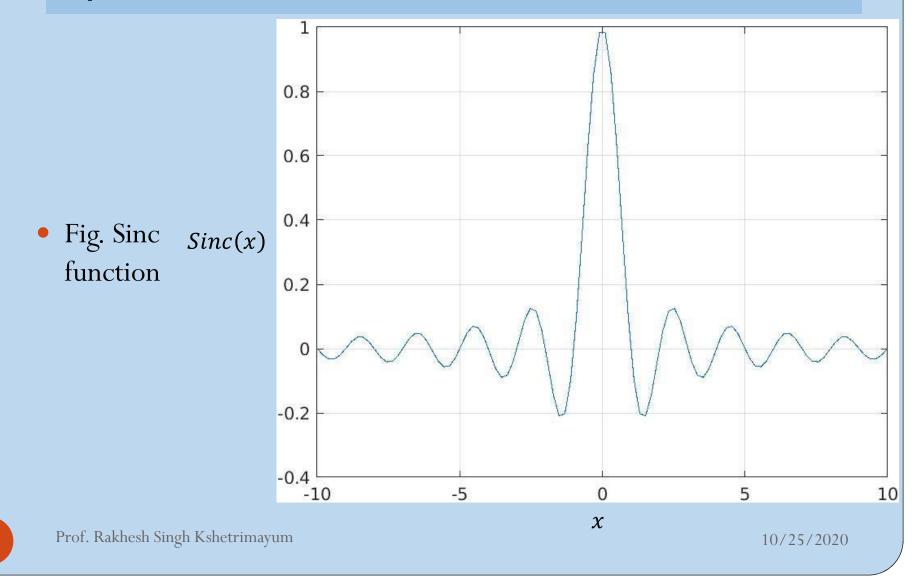
$$= E_0 \int_{-\frac{L_x}{2}}^{\frac{L_x}{2}} e^{(j\beta_x x')} dx' \int_{-\frac{L_y}{2}}^{\frac{L_y}{2}} e^{(j\beta_y y')} dy'$$

$$= E_0 \left\{ \frac{e^{(j\beta_x \frac{L_x}{2})} - e^{(-j\beta_x \frac{L_x}{2})}}{j\beta_x} \right\} \left\{ \frac{e^{(j\beta_y \frac{L_y}{2})} - e^{(-j\beta_y \frac{L_y}{2})}}{j\beta_y} \right\}$$

$$\begin{split} \tilde{E}_{ay} &= 4E_0 \frac{1}{\beta_x} \left\{ \frac{e^{\left(j\beta_x \frac{L_x}{2}\right)} - e^{\left(-j\beta_x \frac{L_x}{2}\right)}}{2j} \right\} \frac{1}{\beta_y} \left\{ \frac{e^{\left(j\beta_y \frac{L_y}{2}\right)} - e^{\left(-j\beta_y \frac{L_y}{2}\right)}}{2j} \right\} \\ &= 4E_0 \left( \frac{\sin\beta_x \frac{L_x}{2}}{\beta_x} \right) \left( \frac{\sin\beta_y \frac{L_y}{2}}{\beta_y} \right) = 4E_0 \frac{L_x}{2} \left( \frac{\sin\beta_x \frac{L_x}{2}}{\beta_x \frac{L_x}{2}} \right) \frac{L_y}{2} \left( \frac{\sin\beta_y \frac{L_y}{2}}{\beta_y \frac{L_y}{2}} \right) \\ &= E_0 L_x L_y sinc \left( \beta_x \frac{L_x}{2} \right) sinc \left( \beta_y \frac{L_y}{2} \right) \end{split}$$

- where  $\beta = \beta_x \hat{x} + \beta_y \hat{y} + \beta_z \hat{z}$  and
- $\beta_x = \beta sin\theta cos\phi$ ,  $\beta_y = \beta sin\theta sin\phi$ ,  $\beta_z = \beta cos\theta$

6



7

• E-plane (electric is maximum along y-axis, hence y-z plane,  $\phi = \frac{\pi}{2}$ )

• 
$$\tilde{E}_{ay} = E_0 L_x L_y sinc\left(\beta_x \frac{L_x}{2}\right) sinc\left(\beta_y \frac{L_y}{2}\right)$$

• where 
$$\beta_x = 0$$
 and  $\beta_y = \beta sin \theta$ 

• Since 
$$sinc\left(\beta_x \frac{L_x}{2}\right) = 1$$
,  $sinc\left(\beta_y \frac{L_y}{2}\right) = sinc\left(\beta sin\theta \frac{L_y}{2}\right)$ , then

• 
$$E_{ff\theta} \cong 2 \frac{j\beta e^{(-j\beta r)}}{4\pi r} [(\tilde{E}_{ay}sin\phi)] =$$
  
 $\frac{j\beta e^{(-j\beta r)}}{2\pi r} [E_0 L_x L_y sinc \left(\beta sin\theta \frac{L_y}{2}\right)]$   
• and  $E_{ff\phi} \cong 2 \frac{j\beta e^{(-j\beta r)}}{4\pi r} [(\tilde{E}_{ay}cos\theta cos\phi)] = 0$ 

• *H*-plane (x-z plane,  $\phi = 0$ )

• 
$$\tilde{E}_{ay} = E_0 L_x L_y sinc\left(\beta_x \frac{L_x}{2}\right) sinc\left(\beta_y \frac{L_y}{2}\right)$$

• where 
$$\beta_x = \beta sin\theta$$
 and  $\beta_y = 0$ 

• 
$$\tilde{E}_{ay} = E_0 L_x L_y sinc\left(\beta sin\theta \frac{L_x}{2}\right)$$

• 
$$E_{ff\theta} \cong 2 \frac{j\beta e^{(-j\beta r)}}{4\pi r} [(\tilde{E}_{ay}sin\phi)] = 0$$
  
and  $E_{ff\phi} \cong 2 \frac{j\beta e^{(-j\beta r)}}{4\pi r} [(\tilde{E}_{ay}cos\theta cos\phi)]$   
 $= \frac{j\beta e^{(-j\beta r)}}{2\pi r} [(E_0 L_x L_y cos\theta sinc(\beta sin\theta \frac{L_x}{2}))]$ 

- Considering E-plane and maximum for  $\theta = 0$ •  $E_{ff\theta} \cong \frac{j\beta e^{(-j\beta r)}}{2\pi r} \left[ E_0 L_x L_y sinc\left(\beta sin\theta \frac{L_y}{2}\right) \right]$ • Nulls occur at  $\beta \frac{L_y}{2} sin \theta_n = n\pi$ ,  $n = 1, 2, \cdots$ • Thus  $sin\theta_n = \frac{n\pi}{\beta \frac{Ly}{2}} = \frac{n\lambda}{L_y}$ •  $BWFN_{E-plane} = 2sin^{-1} \left(\frac{\lambda}{L_{\nu}}\right) rad =$  $114.59sin^{-1}\left(\frac{\lambda}{L_{\nu}}\right)$  degrees
- Similarly,  $BWFN_{H-plane} = 114.59sin^{-1} \left(\frac{\lambda}{L_{r}}\right) degrees$

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- For HPBW,  $sinc^{2}\left(\beta sin\theta \frac{L_{y}}{2}\right) = \frac{1}{2}$ •  $\beta sin\theta \frac{L_{y}}{2} = 1.391 \Rightarrow sin\theta = 1.391 \frac{\lambda}{\pi L_{y}} = \frac{0.443\lambda}{L_{y}}$ •  $HPBW_{E-plane} = 2sin^{-1}\left(\frac{0.443\lambda}{L_{y}}\right)rad$  $= 114.59sin^{-1}\left(\frac{0.443\lambda}{L_{y}}\right)degrees$
- Similarly,
- $HPBW_{H-plane} = 114.59sin^{-1} \left(\frac{0.443\lambda}{L_{\chi}}\right) degrees$
- It is inversely proportional to length in the corresponding cutting plane
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• How do one calculate the directivity (for main beam in broadside direction of the aperture)?

• 
$$D = \frac{4\pi}{\lambda^2} \frac{\left| \iint_{S_a} \vec{E}_a ds' \right|^2}{\iint_{S_a} \left| \vec{E}_a \right|^2 ds'}$$

• For uniform rectangular aperture

$$\vec{E}_{a} = E_{ay}\hat{y} = \begin{cases} E_{0}\hat{y}, & |x| \leq \frac{L_{x}}{2}, |y| \leq \frac{L_{y}}{2}, z = 0\\ 0, & otherwise \end{cases}$$

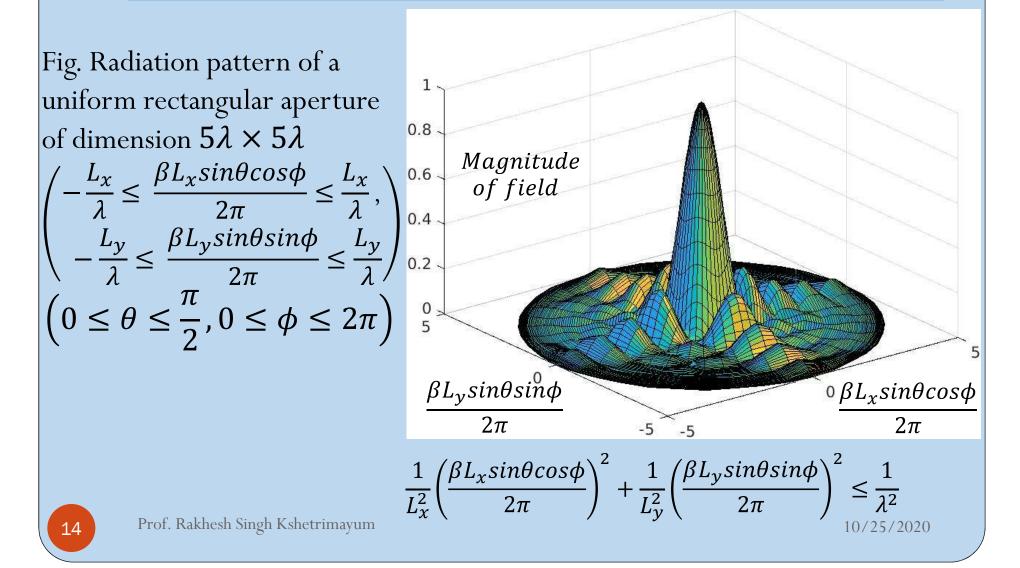
• Hence 
$$D = \frac{4\pi L_x L_y}{\lambda^2}$$

• SLL is greater than -13.26 dB (Proof is given in next slide)

• It is a sinc function, to find the maxima

• 
$$\frac{d}{d}\left(\frac{\sin x}{x}\right) = \frac{x\cos x - \sin x}{x^2} = 0$$

- First maxima solution: x = 0
- is for the main beam
- Second maxima solution:  $x \cong 4.5$
- The value of the second maxima
- $\left|\frac{\sin 4.5}{4.5}\right| = 0.217$
- In dB,
- 20log(0.217) = -13.26dB



• Uniform circular aperture

• 
$$\vec{E}_a = \begin{cases} E_0 \hat{y}, & \rho' \leq a \\ 0, & otherwise \end{cases}$$

- Let us find the 2 D FT
  of the aperture field,
- $\tilde{E}_a(\beta_x, \beta_y)$
- =  $\iint E_a(x', y')e^{(j\beta_x x' + j\beta_y y')}dx'dy'$

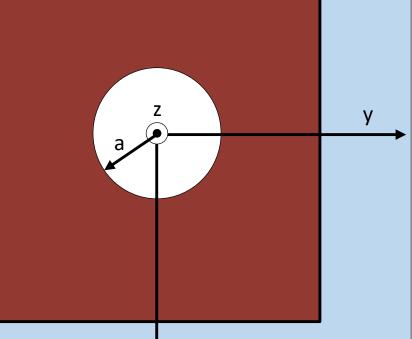


Fig. Circular aperture of radius a in x-y plane

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15

- Better to change to
  - spherical coordinates
- $x' = r' sin\theta' cos\phi'$  and  $y' = r' sin\theta' sin\phi'$
- Since aperture is in x-y plane, we have,  $\theta' = \frac{\pi}{2}$ , hence
- $x' = r' cos \phi'$  and  $y' = r' sin \phi'$
- We know that,
- $\beta_x = \beta sin\theta cos\phi$  and  $\beta_y = \beta sin\theta sin\phi$
- Hence,
- $\beta_x x' + \beta_y y' = \beta r' \sin\theta \cos\phi \cos\phi' + \beta r' \sin\theta \sin\phi \sin\phi'$ =  $\beta r' \sin\theta \cos(\phi - \phi')$

- We also know that area element in x-y plane  $dx'dy' = d\phi'r'dr'$
- Hence
- $\tilde{E}_{a}(\beta sin\theta cos\phi,\beta sin\theta sin\phi) = \int_{0}^{a} \int_{0}^{2\pi} E_{a}(r',\phi') e^{j\beta r'sin\theta cos(\phi-\phi')} d\phi'r'dr'$
- For aperture field distribution with circular symmetry (no field variation w.r.t.  $\phi'$ ),
- $E_a(r',\phi') = E_a(r')$
- Then

• 
$$\tilde{E}_a(\theta,\phi) = \int_0^a \left( \int_0^{2\pi} e^{j\beta r' \sin\theta \cos(\phi - \phi')} d\phi' \right) E_a(r') r' dr'$$

- Note that the inner integral looks like Bessel's integral
- and gives  $2\pi$  times the Bessel's function of first kind
- and of zero order
- Hence,
- $\tilde{E}_a(\theta,\phi) = \int_0^a 2\pi J_0(\beta r' \sin\theta) E_a(r') r' dr'$
- Recall that  $J_0$  is like a cosine function
- but with irregular roots and amplitude taper

- For uniform distribution,
- $E_a(r') = E_0$  along y direction
- $\tilde{E}_{ay}(\theta,\phi) = 2\pi E_0 \int_0^a r' J_0(\beta r' \sin\theta) dr'$
- Using the Bessel's identity
- $\int x J_0(x) \, dx = x J_1(x)$
- Recall that  $J_1$  is like a sine function
  - but with irregular roots and amplitude taper

- Substitute  $\beta r' \sin \theta = \xi$ ,  $d\xi = \beta \sin \theta dr'$ ,
- We have,

• 
$$\tilde{E}_{ay}(\theta,\phi) = 2\pi E_0 \int_0^{\beta a sin\theta} \frac{\xi}{\beta sin\theta} J_0(\xi) \frac{d\xi}{\beta sin\theta}$$
  
=  $2\pi E_0 \frac{\beta a sin\theta J_1(\beta a sin\theta)}{(\beta sin\theta)^2} = 2\pi E_0 \frac{a J_1(\beta a sin\theta)}{\beta sin\theta}$ 

• From (24.1) and (24.2), we have,

• 
$$E_{ff\theta} \cong 2 \frac{j\beta e^{(-j\beta r)}}{4\pi r} [(\tilde{E}_{ay} sin\phi)]$$

• Similarly, 
$$E_{ff\phi} \cong 2 \frac{j\beta e^{(-j\beta r)}}{4\pi r} [(\tilde{E}_{ay} cos\theta cos\phi)]$$

• Hence,

• 
$$E_{ff\theta} \cong \frac{j\beta e^{(-j\beta r)}}{r} \left[ \left( E_0 \frac{aJ_1(\beta asin\theta)}{\beta sin\theta} sin\phi \right) \right]$$
  
• Similarly,  $E_{ff\phi} \cong \frac{j\beta e^{(-j\beta r)}}{r} \left[ \left( E_0 \frac{aJ_1(\beta asin\theta)}{\beta sin\theta} cos\theta cos\phi \right) \right]$   
•  $D = \frac{4\pi}{\lambda^2} \frac{\left| \iint_{S_a} \vec{E}_a ds' \right|^2}{\iint_{S_a} |\vec{E}_a|^2 ds'}$   
•  $\vec{E}_a = \begin{cases} E_0 \hat{y}, \quad \rho' \leq a \\ 0, \quad otherwise \end{cases}$   
• Hence  $D = \frac{4\pi^2 a^2}{\lambda^2}$ 

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• The 3-dB pattern point in E-plane ( $\phi = \frac{\pi}{2}$  or y-z plane) is

• 
$$E_{ff\theta} \cong \frac{j\beta e^{(-j\beta r)}}{r} \left[ \left( E_0 \frac{a J_1(\beta a sin\theta)}{\beta sin\theta} \right) \right]$$

- Similarly,  $E_{ff\phi} \cong 0$
- $\beta asin \theta_1 = 1.6162$

• Or, 
$$sin\theta_1 = \frac{1.6162}{\beta a} = \frac{1.6162\lambda}{\pi 2a} = \frac{1.6162\lambda}{\pi D} = 0.5145 \frac{\lambda}{D}$$

- Hence,  $HPBW_{E-plane} = 2sin^{-1}\left(0.5145\frac{\lambda}{D}\right)rad = 114.59sin^{-1}\left(0.5145\frac{\lambda}{D}\right)degrees$
- where D is diameter

- For BWFN, the first zero or null point is
- $\beta asin \theta_{null} = 3.8317$

• 
$$BWFN_{E-plane} = 2sin^{-1}\left(\frac{3.8317\lambda}{\pi D}\right) = 2sin^{-1}\left(\frac{1.2197\lambda}{D}\right)rad$$
  
=  $114.592sin^{-1}\left(\frac{1.2197\lambda}{D}\right)$ degrees

