

EE540 Advance Electromagnetic Theory & Antennas

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Aperture antennas

- Motivation for Horn antennas:
- *First motivation: Impedance matching*
- Consider a hollow rectangular waveguide operating in dominant TE_{10} mode
- What is the wave impedance?
- $Z_{TE} = \frac{\eta_0}{\sqrt{1 - \left(\frac{f_c}{f}\right)^2}}$
- What is cut-off frequency for the dominant TE_{10} mode inside rectangular waveguide?

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- $f_c = \frac{1}{2a\sqrt{\mu\epsilon}} = \frac{c}{2a}$
- For X-band waveguide with $a=22.86$ mm and $b=10.16$ mm
- $f_c = 6.56$ GHz
- Therefore around 10 GHz,
- $Z_{TE} = 499 \Omega$
- But the wave impedance for free space is around 377Ω
- So reflection coefficient is
- $\Gamma = \frac{377-499}{377+499} = -.139$

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- If we increase the a five times, then
- $a=114.3\text{mm}$
- What is new cut-off frequency for the dominant TE_{10} modes inside this new rectangular waveguide?
- $f_c = \frac{1}{2a\sqrt{\mu\epsilon}} = \frac{c}{2a}$
- $f_c = 1.311 \text{ GHz}$
- Therefore around 10 GHz, $Z_{TE} = 380 \Omega$
- Hence, reflection coefficient for this case is
- $\Gamma = \frac{377-380}{377+380} = -.00396$
- Impedance matching is much better in this case

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- *Second motivation:*
- Directivity enhancements:
- Directivity of open ended waveguide is given by
- $D \cong \frac{4\pi}{\lambda^2} (0.81ab)$
- What if we increase a and b (dimensions of the rectangular waveguide)?
- Say, if we double both a and b,
 - directivity will increase fourfold

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- *Third motivation:*
- Narrower beam:

E-plane

$$E_{ff\theta} \cong \frac{j\beta e^{(-j\beta r)}}{r} \left[\left(abE_0 \frac{1}{\pi^2} \text{sinc} \left(\beta \frac{b}{2} \sin\theta \right) \right) \right]$$

- For example, E-plane has sinc function and its argument
- $\beta \frac{b}{2} \sin\theta$
- What will happen b becomes five times longer?

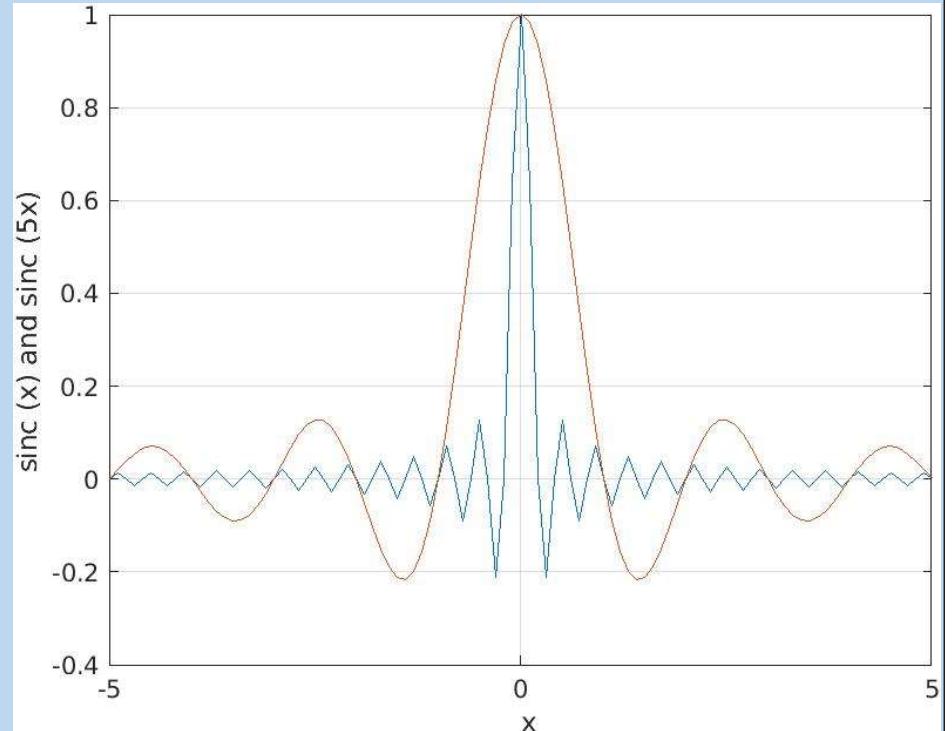
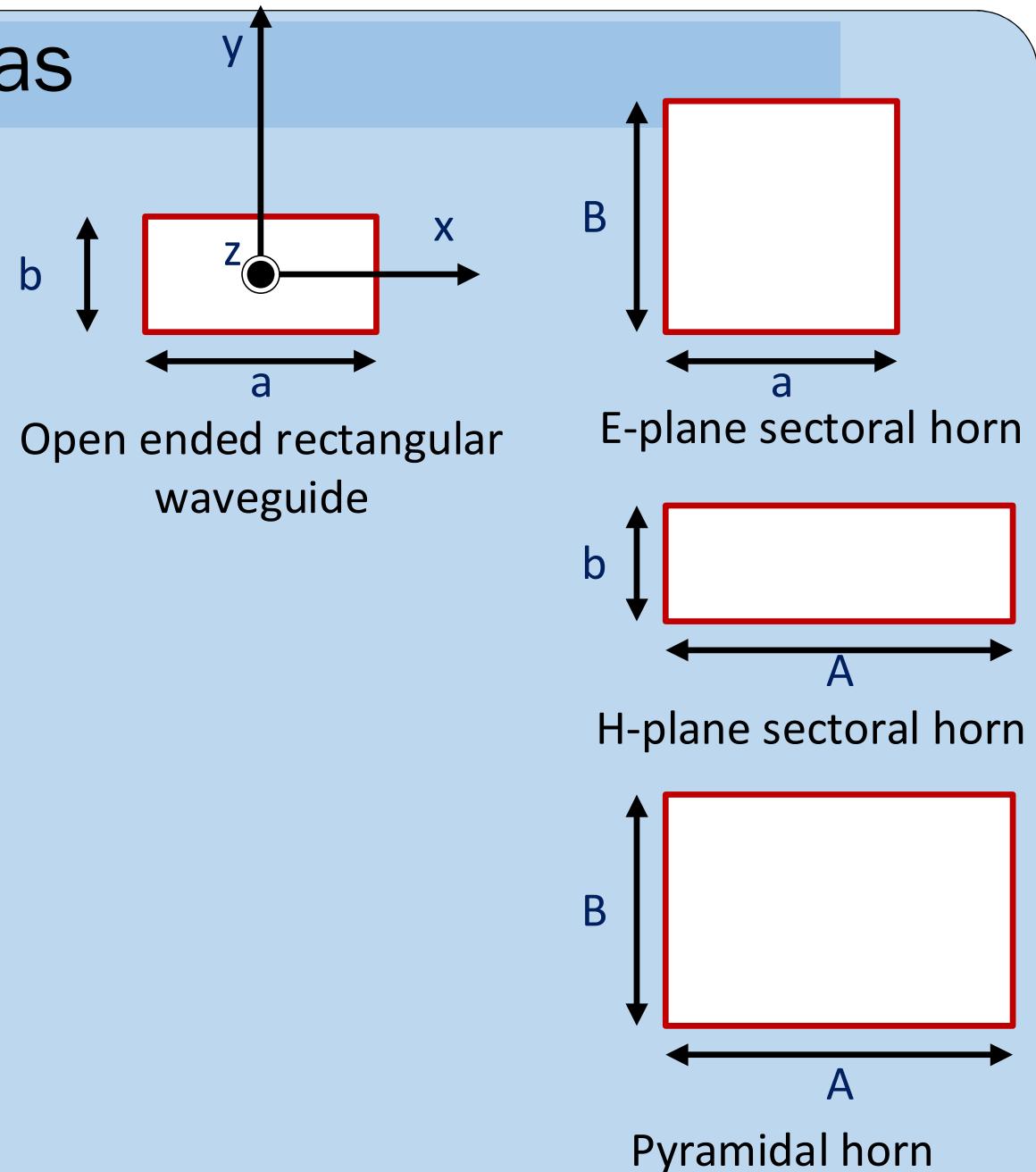


Fig. Sinc (5x) and Sinc (x) vs x

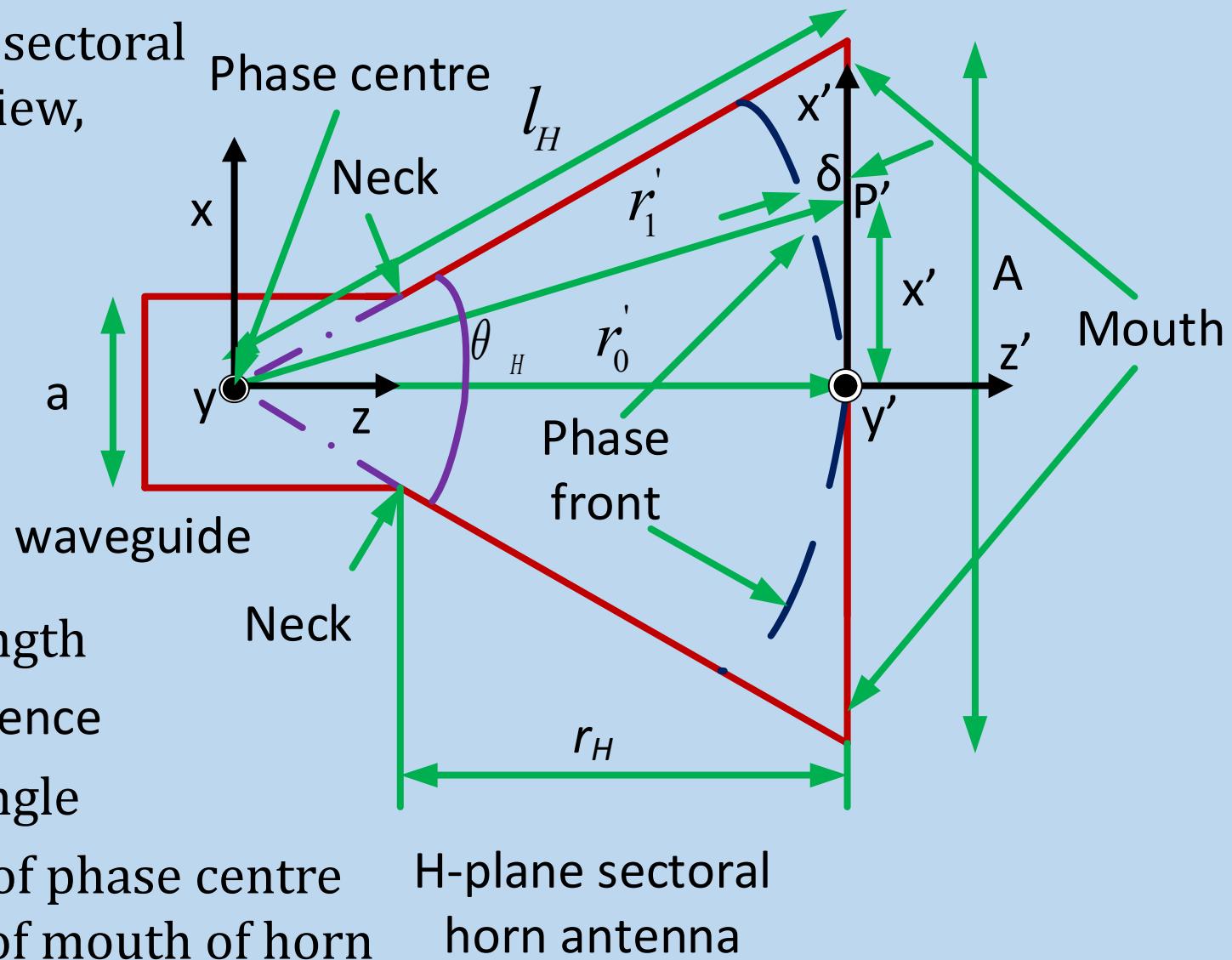
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- Fig. Open ended waveguide to horn antenna (front view)



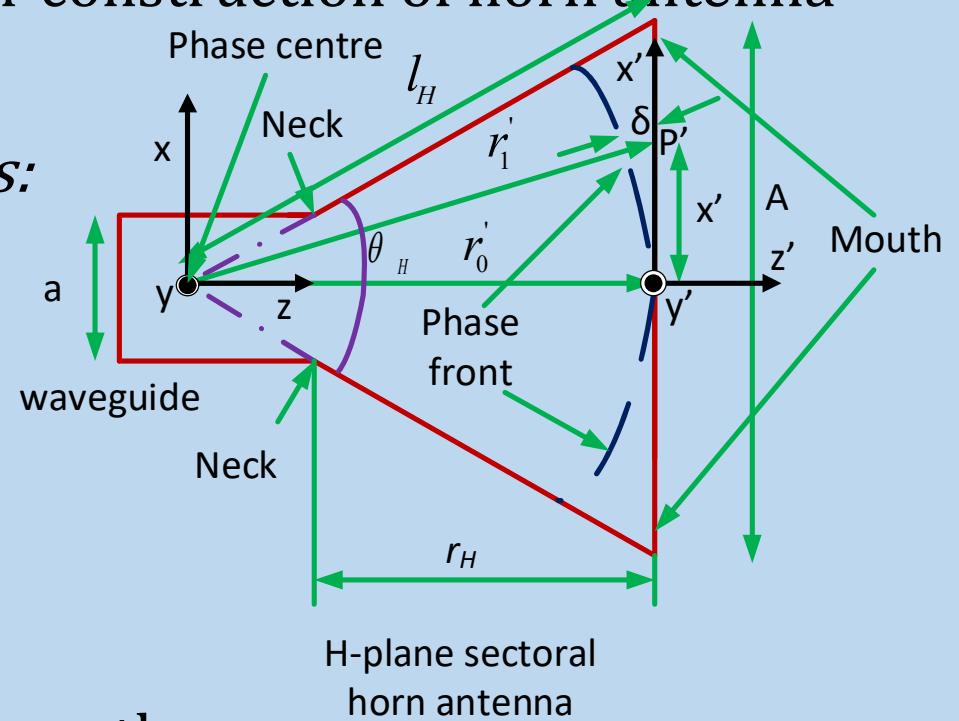
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- Fig. H-plane sectoral horn (side view, x-z cut)



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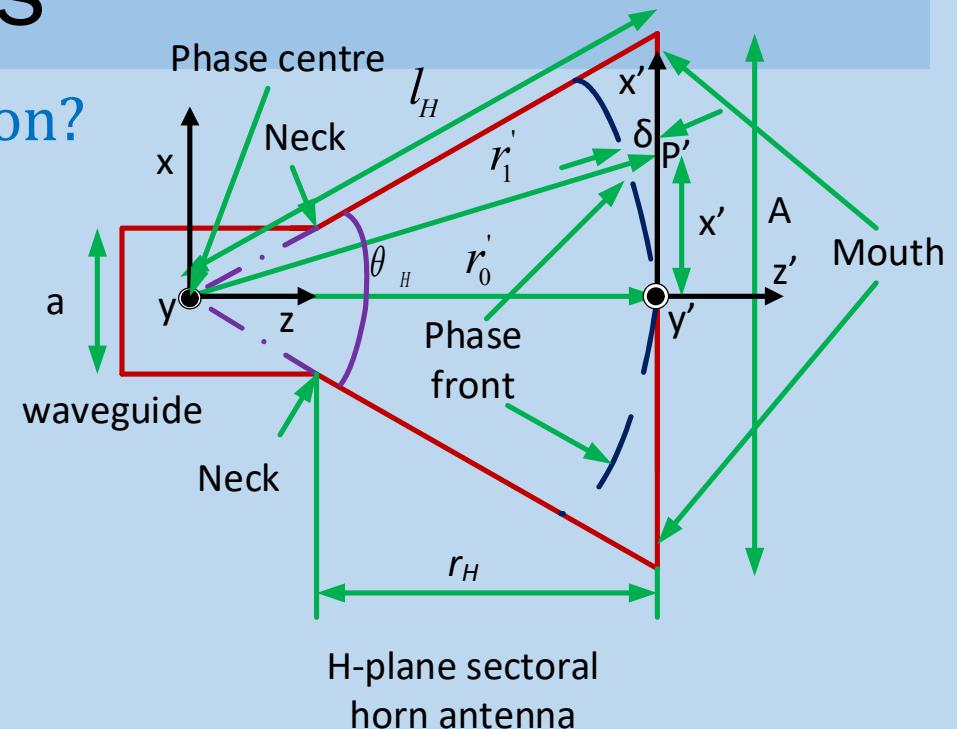
- Two dimensions required for construction of horn antenna
- A and r_H
- *Simple geometrical relations:*
- Flare angle
- $\theta_H = 2\tan^{-1} \left(\frac{A}{2r'_0} \right)$
- Flaring length
- $l_H^2 = \frac{A^2}{4} + (r'_0)^2$
- Distance between neck and mouth r_H
- $r_H = (A - a) \sqrt{\left(\frac{l_H}{A}\right)^2 - \frac{1}{4}}$



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- How do you get last equation?
- Note that
- $\tan\left(\frac{\theta_H}{2}\right) = \frac{A}{2r'_0} = \frac{a}{2(r'_0 - r_H)}$
- Inverting and cancelling 2
- $\frac{r'_0}{A} = \frac{(r'_0 - r_H)}{a}$
- Cross multiplying
- $ar'_0 = Ar'_0 - Ar_H \Rightarrow Ar_H = (A - a)r'_0 \Rightarrow r_H = (A - a)\frac{r'_0}{A}$

$$\Rightarrow r_H = (A - a)\frac{\sqrt{l_H^2 - \frac{A^2}{4}}}{A} = (A - a)\sqrt{\left(\frac{l_H}{A}\right)^2 - \frac{1}{4}}$$



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- Note that at any point P',
 - since the distance r'_0 and r'_1 are not same
- Hence there will be some phase change
- *How much is the phase change?*
- Let us find the difference in the distance $r'_1 - r'_0$ first
- Applying the Pythagoras theorem,

- $r'_1 = \sqrt{(r'_0)^2 + (x')^2} = r'_0 \left(1 + \left(\frac{x'}{r'_0} \right)^2 \right)^{1/2} \cong r'_0 \left(1 + \frac{1}{2} \left(\frac{x'}{r'_0} \right)^2 \right)$
- Hence, $r'_1 - r'_0 \cong \frac{1}{2r'_0} (x')^2$

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- $\delta = r'_1 - r'_0 \cong \frac{1}{2r'_0} (x')^2$
- We can observe that $r'_1 - r'_0$ has quadratic variation with x'
- Hence phase has quadratic variation
- *First step:* so we can write the aperture field variation as
- $E_{ay} = E_0 \cos\left(\frac{\pi x'}{A}\right) e^{-j\beta \frac{1}{2r'_0} (x')^2}$ inside the aperture and zero elsewhere
 - $E_0 \cos\left(\frac{\pi x'}{A}\right)$ is the amplitude variation and
 - $e^{-j\beta \frac{1}{2r'_0} (x')^2}$ is the phase variation

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Table: Amplitude and phase variation of H-sectoral horn antenna

| Aperture field parameter | Expression | $x' = 0$ (centre of mouth) | $x' = -A/2$ (extreme edge of mouth) | $x' = A/2$ (extreme edge of mouth) |
|--------------------------|---|-------------------------------|--|--|
| Amplitude | $E_0 \cos\left(\frac{\pi x'}{A}\right)$ | E_0 | 0 | 0 |
| Phase | $e^{-j\beta \frac{1}{2r'_0} (x')^2}$ | 1 | $e^{-j\beta \frac{1}{8r'_0} (A)^2}$ (maximum) | $e^{-j\beta \frac{1}{8r'_0} (A)^2}$ (maximum) |

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- Second step:
- Take the 2-D FT of aperture field

- $\tilde{E}_{ay} = \int_{-A/2}^{A/2} E_0 \cos\left(\frac{\pi x'}{A}\right) e^{-j\beta \frac{1}{2r'_0}(x')^2} e^{j\beta_x x'} dx' \int_{-b/2}^{b/2} e^{j\beta_y y'} dy'$
- $= E_0 \left\{ \frac{1}{2} \sqrt{\frac{\pi r'_0}{\beta}} I(\theta, \phi) \right\} \left\{ b sinc\left(\beta_y \frac{b}{2}\right) \right\}$
- where $I(\theta, \phi) = A_1(\beta_{x1})(f(s'_2) - f(s'_1)) + A_2(\beta_{x2})(f(t'_2) - f(t'_1))$
- $A_1(\beta_{x1}) = e^{j\frac{r'_0(\beta_{x1})^2}{2\beta}}, A_2(\beta_{x2}) = e^{j\frac{r'_0(\beta_{x2})^2}{2\beta}},$
- $\beta_{x1} = \beta_x + \frac{\pi}{A}, \beta_{x2} = \beta_x - \frac{\pi}{A}$

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- $f(\xi') = C(\xi') - jS(\xi')$,
- $C(\xi') = \int_0^{\xi'} \cos\left(\frac{\pi}{2}\tau^2\right) d\tau, S(\xi') = \int_0^{\xi'} \sin\left(\frac{\pi}{2}\tau^2\right) d\tau$
- $s'_1 = \frac{1}{\sqrt{\pi\beta r'_0}} \left(-\frac{\beta A}{2} - \beta_{x1} r'_0 \right)$
- $s'_2 = \frac{1}{\sqrt{\pi\beta r'_0}} \left(+\frac{\beta A}{2} - \beta_{x1} r'_0 \right)$
- $t'_1 = \frac{1}{\sqrt{\pi\beta r'_0}} \left(-\frac{\beta A}{2} - \beta_{x2} r'_0 \right)$
- $t'_2 = \frac{1}{\sqrt{\pi\beta r'_0}} \left(+\frac{\beta A}{2} - \beta_{x2} r'_0 \right)$

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- *Third step:*
- Find the fields using equations (24.1) and (24.2)
- Recall for TE₁₀ mode waveguide,
 - $E_{ay} = E_0 \cos\left(\frac{\pi x'}{A}\right) e^{-j\beta \frac{1}{2r_0}(x')^2}$ inside the aperture and zero elsewhere
 - Note that there is no infinite ground plane
 - Besides, inherently, waveguide has both aperture electric and magnetic fields
 - And their relation is $H_{ax} = -\frac{E_{ay}}{Z_g} \Rightarrow E_{ay} = -H_{ax}Z_g$

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- Revisiting equations (24.1) and (24.2) for this aperture antenna,
- $E_{ff\theta} \cong \frac{j\beta e^{(-j\beta r)}}{4\pi r} [\eta(\tilde{H}_{ay} \cos\theta \cos\phi - \tilde{H}_{ax} \cos\theta \sin\phi) + (\tilde{E}_{ay} \sin\phi + \tilde{E}_{ax} \cos\phi)]$ ---(24.1)
- $E_{ff\phi} \cong \frac{j\beta e^{(-j\beta r)}}{4\pi r} [-\eta(\tilde{H}_{ay} \sin\phi + \tilde{H}_{ax} \cos\phi) + (\tilde{E}_{ay} \cos\theta \cos\phi - \tilde{E}_{ax} \cos\theta \sin\phi)]$ ---(24.2)
- Since there is contribution from \tilde{E}_{ay} and \tilde{H}_{ax} , it is simplified as
- $E_{ff\theta} \cong \frac{j\beta e^{(-j\beta r)}}{4\pi r} [\eta(-\tilde{H}_{ax} \cos\theta \sin\phi) + (\tilde{E}_{ay} \sin\phi)]$ ---(24.1)
- $E_{ff\phi} \cong \frac{j\beta e^{(-j\beta r)}}{4\pi r} [-\eta(\tilde{H}_{ax} \cos\phi) + (\tilde{E}_{ay} \cos\theta \cos\phi)]$ ---(24.2)

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- Using the relation $E_{ay} = -H_{ax}Z_g$ (note that η for the waveguide case should be the wave impedance for the mode we are considering), it can be simplified as
- $E_{ff\theta} \cong \frac{j\beta e^{(-j\beta r)}}{4\pi r} [(\tilde{E}_{ay} \cos\theta \sin\phi) + (\tilde{E}_{ay} \sin\phi)] \quad \text{---(24.1)}$
- $E_{ff\phi} \cong \frac{j\beta e^{(-j\beta r)}}{4\pi r} [(\tilde{E}_{ay} \cos\phi) + (\tilde{E}_{ay} \cos\theta \cos\phi)] \quad \text{---(24.2)}$
- Further simplification after combining common terms,
- $E_{ff\theta} \cong \frac{j\beta e^{(-j\beta r)}}{4\pi r} (1 + \cos\theta) [\tilde{E}_{ay} \sin\phi] \quad \text{---(27.1)}$
- $E_{ff\phi} \cong \frac{j\beta e^{(-j\beta r)}}{4\pi r} (1 + \cos\theta) [\tilde{E}_{ay} \cos\phi] \quad \text{---(27.2)}$

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- where $\tilde{E}_{ay} = E_0 \left\{ \frac{1}{2} \sqrt{\frac{\pi r'_0}{\beta}} I(\theta, \phi) \right\} \left\{ b s i n c \left(\beta_y \frac{b}{2} \right) \right\}$
- $I(\theta, \phi) = A_1(\beta_{x1})(f(s'_2) - f(s'_1)) + A_2(\beta_{x2})(f(t'_2) - f(t'_1))$
- E-plane patterns ($\phi = \frac{\pi}{2}$):
 - $E_{ff\theta} \cong \frac{j\beta e^{(-j\beta r)}}{4\pi r} (1 + \cos\theta) [\tilde{E}_{ay}]$ and $E_{ff\phi} \cong 0$
 - $\beta_x = \beta \sin\theta \cos\phi = 0$ and $\beta_y = \beta \sin\theta \sin\phi = \beta \sin\theta$
- $\beta_{x1} = \frac{\pi}{A}, \beta_{x2} = -\frac{\pi}{A}, A_1 \left(\frac{\pi}{A} \right) = e^{j \frac{r'_0 (\frac{\pi}{A})^2}{2\beta}}, A_2 \left(-\frac{\pi}{A} \right) = e^{j \frac{r'_0 (\frac{\pi}{A})^2}{2\beta}}$
- Hence $A_1 \left(\frac{\pi}{A} \right) = A_2 \left(-\frac{\pi}{A} \right)$

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- $\therefore I(\theta, \phi) = A_1 \left(\frac{\pi}{A} \right) (f(s'_2) - f(s'_1)) + (f(t'_2) - f(t'_1))$
- $s'_1 = \frac{1}{\sqrt{\pi\beta r'_0}} \left(-\frac{\beta A}{2} - \frac{\pi}{A} r'_0 \right)$ and $s'_2 = \frac{1}{\sqrt{\pi\beta r'_0}} \left(+\frac{\beta A}{2} - \frac{\pi}{A} r'_0 \right)$
- $t'_1 = \frac{1}{\sqrt{\pi\beta r'_0}} \left(-\frac{\beta A}{2} + \frac{\pi}{A} r'_0 \right)$ and $t'_2 = \frac{1}{\sqrt{\pi\beta r'_0}} \left(+\frac{\beta A}{2} + \frac{\pi}{A} r'_0 \right)$
- Hence $s'_2 = -t'_1$ and $s'_1 = -t'_2$
- $f(\xi') = C(\xi') - jS(\xi')$ where $C(-\xi') = -C(\xi')$ and $S(-\xi') = -S(\xi')$ and $C(\xi')$ and $S(\xi')$ are sine and cosine Fresnel integrals
- Therefore, $f(-\xi') = -f(\xi')$
- $\therefore I(\theta, \phi) = A_1 \left(\frac{\pi}{A} \right) 2(f(s'_2) - f(s'_1))$

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- Hence
- $E_{ff\theta} \cong \frac{j\beta b E_0 e^{(-j\beta r)}}{4\pi r} \sqrt{\frac{\pi r'_0}{\beta}} \left[A_1 \left(\frac{\pi}{A} \right) 2(f(s'_2) - f(s'_1)) \right] \left(\frac{1+\cos\theta}{2} \right) \left\{ \text{sinc} \left(\beta \sin\theta \frac{b}{2} \right) \right\}$
- H-plane patterns ($\phi = 0$):
- $E_{ff\theta} \cong 0$ and $E_{ff\phi} \cong \frac{j\beta e^{(-j\beta r)}}{4\pi r} (1 + \cos\theta) [\tilde{E}_{ay}]$
- $\beta_x = \beta \sin\theta \cos\phi = \beta \sin\theta$ and $\beta_y = \beta \sin\theta \sin\phi = 0$
- $\beta_{x1} = \beta_x + \frac{\pi}{A}$, $\beta_{x2} = \beta_x - \frac{\pi}{A}$, $A_1(\beta_{x1}) = e^{j\frac{r'_0(\beta_{x1})^2}{2\beta}}$, $A_2(\beta_{x2}) = e^{j\frac{r'_0(\beta_{x2})^2}{2\beta}}$

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- $\tilde{E}_{ay} = E_0 \left\{ \frac{1}{2} \sqrt{\frac{\pi r'_0}{\beta}} I(\theta, \phi) \right\} \left\{ b s i n c \left(\beta_y \frac{b}{2} \right) \right\}$
- $I(\theta, \phi) = A_1(\beta_{x1})(f(s'_2) - f(s'_1)) + A_2(\beta_{x2})(f(t'_2) - f(t'_1))$
- Hence
- $E_{ff\phi} \cong \frac{j\beta b E_0 e^{(-j\beta r)}}{4\pi r} \sqrt{\frac{\pi r'_0}{\beta}} [A_1(\beta_{x1})(f(s'_2) - f(s'_1)) + A_2(\beta_{x2})(f(t'_2) - f(t'_1))] \left(\frac{1+cos\theta}{2} \right)$
- Directivity
- $D_H = \frac{4\pi}{\lambda^2} \epsilon_t \epsilon_{ph}^H A b$

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- $\epsilon_t = \frac{8}{\pi^2} \cong 0.81$
- $\epsilon_{ph}^H = \frac{\pi^2}{64t} \left[(C(p_1) - C(p_2))^2 + (S(p_1) - S(p_2))^2 \right]$

How do we find t for maximum phase error?

- Maximum phase error is for $x' = \frac{A}{2}$,
- hence $\beta \frac{1}{2r'_0} (x')^2 = \beta \frac{1}{8r'_0} (A)^2 = 2\pi \frac{A^2}{8\lambda r'_0}$
- Let us equate the maximum phase error to $2\pi t$
- So, $\frac{A^2}{8\lambda r'_0} = t$
-

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- What are parameters p_1 and p_2 ?
- $s'_1 = \frac{1}{\sqrt{\pi\beta r'_0}} \left(-\frac{\beta A}{2} - \beta_{x1} r'_0 \right) = \frac{1}{\sqrt{\pi\beta r'_0}} \left(-\frac{\beta A}{2} - \left(\beta \sin\theta \cos\phi + \frac{\pi}{A} \right) r'_0 \right) = \frac{1}{\sqrt{\pi\beta r'_0}} \left(-\frac{\beta A}{2} - \left(\beta u + \frac{\pi}{A} \right) r'_0 \right)$
- For $u = \sin\theta \cos\phi = 0$,
- $p_1 = -s'_1 = \frac{1}{\sqrt{\pi\beta r'_0}} \left(\frac{\beta A}{2} + \left(\frac{\pi}{A} \right) r'_0 \right) = t'_2$
- $p_2 = s'_2 = \frac{1}{\sqrt{\pi\beta r'_0}} \left(+\frac{\beta A}{2} - \left(\frac{\pi}{A} \right) r'_0 \right) = -t'_1$

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- Can we express p_1 and p_2 in terms of t ($\sqrt{8\lambda r'_0 t} = A$)?

$$\bullet p_1 = \frac{1}{\sqrt{\pi \beta r'_0}} \left(\frac{\beta \sqrt{8\lambda r'_0 t}}{2} + \left(\frac{\pi}{\sqrt{8\lambda r'_0 t}} \right) r'_0 \right) = 2\sqrt{t} \left(1 + \frac{1}{8t} \right)$$

$$\bullet p_2 = s'_2 = \frac{1}{\sqrt{\pi \beta r'_0}} \left(+ \frac{\beta A}{2} - \left(\frac{\pi}{A} \right) r'_0 \right) = 2\sqrt{t} \left(-1 + \frac{1}{8t} \right)$$

- Is directivity dependent on A ?

- Yes and optimal value is for $A = \sqrt{3\lambda r'_0}$

- What is the value of t for this? $t = \frac{A^2}{8\lambda r'_0} = \frac{3\lambda r'_0}{8\lambda r'_0} = \frac{3}{8}$

- $D_H = \frac{4\pi}{\lambda^2} \epsilon_t \epsilon_{ph}^H A b$