EE540 Advance Electromagnetic Theory & Antennas

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भारतीय प्रौद्योगिकी संस्थान गुवाहाटी INDIAN INSTITUTE OF TECHNOLOGY GUWAHATI Time Harmonic Electromagnetic Fields

- For time harmonic fields:
- Separate the space and time dependence as $\vec{F}(x, y, z, t) = \vec{F}(x, y, z)e^{j\omega t}$
 - where $\boldsymbol{\omega}$ is the angular frequency of the time varying field
- Time derivative time harmonic fields

$$\frac{\partial \vec{F}(x,y,z,t)}{\partial t} = \vec{F}(x,y,z)\frac{\partial e^{j\omega t}}{\partial t} = \vec{F}(x,y,z)j\omega e^{j\omega t} = j\omega\vec{F}(x,y,z)e^{j\omega t} = j\omega\vec{F}(x,y,z,t)$$

• Integration of time harmonic field w.r.t. time

$$\int \vec{F}(x,y,z,t)dt = \int \vec{F}(x,y,z)e^{j\omega t}dt = \vec{F}(x,y,z)\int e^{j\omega t}dt = \vec{F}(x,y,z)\frac{e^{j\omega t}}{j\omega} + c_1 = \frac{\vec{F}(x,y,z,t)}{j\omega} + c_1$$

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Time Harmonic Electromagnetic Fields

- where c_1 is a constant of integration
- Since we are interested only in time varying quantities, we can take, c₁=0
- So we can define

$$\frac{\partial}{\partial t} \equiv j\omega; \frac{\partial^2}{\partial t^2} = j\omega \times j\omega = -\omega^2$$

- For time harmonic fields, we can modify Maxwell's equations
 - By suppressing time dependence completely
 - It is called Maxwell's equations in phasor form

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- Maxwell's equations in phasor form
 - all time dependence has been suppressed and replace

$$\nabla \bullet \widetilde{D} = \rho_{\iota}$$

$$\nabla \bullet \widetilde{B} = 0$$

$$\nabla \times \widetilde{H} = j\omega \widetilde{D} + \widetilde{J}$$

$$\nabla \times \tilde{E} = -j\omega \tilde{B}$$

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 $\frac{\partial}{\partial t} \equiv j\omega$

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Vector Potentials

- Usually for antennas, we find fields from the current density
- Observation: difficult to find fields directly from current density
 - calculations are highly complex and tedious
- We know that it is easier to
 - find electric field from electric potential than
 - directly finding electric field (Why easier?) $\vec{E} = -\nabla V$
- A major simplification is possible when
 - we find the magnetic vector potential first from current density (Why easier?)
 - and find the fields from it

$$\vec{A}(\vec{r}) = \frac{\mu}{4\pi} \int_{V} \vec{J}(\vec{r}') \frac{e^{-j\beta |\vec{r} - \vec{r}'|}}{|\vec{r} - \vec{r}'|} dv'$$

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Vector Potentials

- One of the Maxwell's divergence equation $\nabla \bullet \vec{B} = 0$
- Hence, we can write

$$\therefore \nabla \bullet \vec{B} = 0, \nabla \bullet \left(\nabla \times \vec{A} \right) = 0 \therefore \vec{B} = \nabla \times \vec{A}$$

- It means that we can find magnetic flux density
 - from the curl of magnetic vector potential
- Putting this in the following Maxwell's curl equation

$$\Rightarrow \nabla \times \vec{E} = -\frac{\partial \left(\nabla \times \vec{A}\right)}{\partial t}$$

$$\nabla \times \vec{E} = -\frac{\partial B}{\partial t}$$

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Vector Potentials

- which can be rewritten as • From vector analysis $\nabla \times \left(\vec{E} + \frac{\partial \vec{A}}{\partial t}\right) = 0 \Rightarrow \nabla \times \left(\nabla V\right) = 0 \Rightarrow \nabla \times \left(-\nabla V\right) = 0$ • For time varying fields, $\vec{E} = -\nabla V - \frac{\partial \vec{A}}{\partial t}$
- For electrostatics, \bar{E}

$$\vec{E} = -\nabla V$$

- The additional term for time varying fields is $-\frac{\partial A}{\partial t}$
- It means that we can find electric field
 - from the negative of the gradient of electric potential
 - and the negative of the time derivative of magnetic vector potential

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