# EE540 Advance Electromagnetic Theory & Antennas

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#### Wave equations from Maxwell's equations

- In linear isotropic medium,
- the two Maxwell curl equations for time harmonic fields are

 $\nabla \times \vec{E} = -j\omega\mu\vec{H} \quad \text{(Equation 1)}$  $\nabla \times \vec{H} = j\omega\varepsilon\vec{E} + \vec{J} = j\omega\varepsilon\vec{E} + \sigma\vec{E} = (j\omega\varepsilon + \sigma)\vec{E} \quad \text{(Equation 2)}$ 

- To solve for electric fields,
  - we can take curl on the first equation, for LHS

$$\nabla \times \nabla \times \vec{E} = \nabla \times \nabla \times \vec{E} = \nabla (\nabla \bullet \vec{E}) - \nabla^2 \vec{E}$$

- For a charge free region, wave propagation in free space,  $\nabla \bullet \vec{E} = 0$ 
  - therefore  $\nabla \times \nabla \times \vec{E} = -\nabla^2 \vec{E}$



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- eliminate the curl of magnetic field vector
- in the RHS, after taking curl of the first equation
  - by using the second equation given above

$$\therefore \nabla \times \nabla \times \vec{E} = -\nabla^2 \vec{E} = -j\omega\mu(\nabla \times \vec{H}) = -j\omega\mu(j\omega\varepsilon + \sigma)\vec{E} = -\gamma^2 \vec{E}$$

• Hence,

$$\nabla^2 \vec{E} - \gamma^2 \vec{E} = 0$$

• Similarly,

 $\nabla^2 \vec{H} - \gamma^2 \vec{H} = 0$ 

- Note that Helmholtz wave equation is a vector wave equation
  - So Laplacian operator  $\nabla^2$  will operate on each components of vector  $\vec{E}/\vec{H}$



- Solution of wave equation in Cartesian coordinates  $\nabla^2 \vec{E} \gamma^2 \vec{E} = 0$ 
  - Electric field is a vector and will have 3 components

 $\vec{E} = E_x \hat{x} + E_y \hat{y} + E_z \hat{z}$ 

- Laplacian will operate on all three components
- Consider Laplacian operating on  $E_{\chi}$

$$\nabla^{2} E_{x} - \gamma^{2} E_{x} = \left(\frac{\partial^{2}}{\partial x^{2}} + \frac{\partial^{2}}{\partial y^{2}} + \frac{\partial^{2}}{\partial z^{2}}\right) E_{x} - \gamma^{2} E_{x} = 0$$

- Applying method of separation of variables, assume  $E_x = X(x)Y(y)Z(z)$
- Rewrite wave equation as
- Divide by XYZ on both sides, we have

$$\left(YZ\frac{\partial^2 X}{\partial x^2} + XZ\frac{\partial^2 Y}{\partial y^2} + XY\frac{\partial^2 Z}{\partial z^2}\right) = \gamma^2 XYZ$$
$$\left(\frac{1}{X}\frac{\partial^2 X}{\partial x^2} + \frac{1}{Y}\frac{\partial^2 Y}{\partial y^2} + \frac{1}{Z}\frac{\partial^2 Z}{\partial z^2}\right) = \gamma^2$$

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- RHS is a constant, LHS first, second and third term
  - are totally dependent on x, y and z respectively

Hence assume 
$$\frac{1}{X}\frac{\partial^2 X}{\partial x^2} = \gamma_x^2, \frac{1}{Y}\frac{\partial^2 Y}{\partial y^2} = \gamma_y^2, \frac{1}{Z}\frac{\partial^2 Z}{\partial z^2} = \gamma$$

$$X = C_x e^{\pm \gamma_x x}, Y = C_y e^{\pm \gamma_y y}, Z = C_z e^{\pm \gamma_z z}$$

• Therefore,

whose solutions are

$$E_x = XYZ = C_x C_y C_z e^{\pm \gamma_x x} e^{\pm \gamma_y y} e^{\pm \gamma_z z} = C_1 e^{\pm \gamma_x x \pm \gamma_y y \pm \gamma_z z}$$

• Similarly  $E_y = C_2 e^{\pm \gamma_x x \pm \gamma_y y \pm \gamma_z z}, E_z = C_3 e^{\pm \gamma_x x \pm \gamma_y y \pm \gamma_z z}$ 

• Therefore,

$$\vec{E} = E_x \hat{x} + E_y \hat{y} + E_z \hat{z} = (C_1 \hat{x} + C_2 \hat{y} + C_3 \hat{z}) e^{\pm (\gamma_x x + \gamma_y y + \gamma_z z)} = \vec{C} e^{\pm (\gamma_x x + \gamma_y y + \gamma_z z)}$$

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5



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• For

Reduce

 $\vec{\gamma} = \gamma_x \hat{x} + \gamma_y \hat{y} + \gamma_z \hat{z}, \vec{r} = x\hat{x} + y\hat{y} + z\hat{z}$  $\vec{E} = \vec{C}e^{\pm(\vec{\gamma} \cdot \vec{r})}$ 

• For time harmonic fields

$$\vec{E} = \vec{C}e^{j\omega t \pm (\vec{\gamma} \bullet \vec{r})} = \vec{C}e^{j\omega t \pm ((\vec{\alpha} + j\vec{\beta}) \bullet \vec{r})} = \vec{C}e^{\pm \vec{\alpha} \bullet \vec{r}}e^{j(\omega t \pm \vec{\beta} \bullet \vec{r})}$$

• Taking real parts (plane wave)

$$\operatorname{Re}\left\{\vec{E}\right\} = \operatorname{Re}\left\{\vec{C}e^{\pm\vec{\alpha}\cdot\vec{r}}e^{j\left(\omega t \pm \vec{\beta}\cdot\vec{r}\right)}\right\} = \vec{C}e^{\pm\vec{\alpha}\cdot\vec{r}}\cos\left(\omega t \pm \vec{\beta}\cdot\vec{r}\right)$$

- Simplifications:
- Assume x-polarized (electric field has only x-component,  $\vec{E} = E_x \hat{x}$ )
- wave propagating in positive z-direction,  $\vec{\gamma} = \gamma_z \hat{z} = (\alpha_z + j\beta_z)\hat{z}$

$$\operatorname{Re}\left\{\vec{E}\right\} = \operatorname{Re}\left\{C\hat{x}e^{-\alpha_{z}z}e^{j(\omega t - \beta_{z}z)}\right\} = Ce^{-\alpha_{z}z}\cos(\omega t - \beta_{z}z)\hat{x}$$

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6



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- General Orthogonal Coordinate System
  - Introduce a new set of coordinates
  - (extension of Cartesian coordinates (x,y,z))
    - $a_1 = a_1(x, y, z), a_2 = a_2(x, y, z) \text{ and } a_3 = a_3(x, y, z)$
  - where the directions at any point indicated by  $\widehat{a}_1$ ,  $\widehat{a}_2$  and  $\widehat{a}_3$ 
    - are orthogonal (perpendicular) to each other
  - is referred to as a set of orthogonal curvilinear coordinates
  - For example,
  - Cylindrical coordinate system
  - $a_1 = \rho, a_2 = \varphi, a_3 = z$
  - $\hat{a}_1 = \hat{\rho}$ ,  $\hat{a}_2 = \hat{\varphi}$ ,  $\hat{a}_3 = \hat{z}$
  - Relation between Cartesian and Cylindrical coordinates
  - $x = \rho \cos\varphi$ ,  $y = \rho \sin\varphi$ , z = z



- With each coordinate is associated a scale factor
  - s<sub>1</sub>, s<sub>2</sub> and s<sub>3</sub> respectively
  - The scale factor gives a measure of
  - how a change in the coordinate changes the position of a point
  - Calculation of scale factors
  - Relation between Cartesian and Cylindrical coordinates

• 
$$x = \rho \cos\varphi, y = \rho \sin\varphi, z = z$$
  
•  $s_1 = \sqrt{\left(\frac{\partial x}{\partial \rho}\right)^2 + \left(\frac{\partial y}{\partial \rho}\right)^2 + \left(\frac{\partial z}{\partial \rho}\right)^2} = \sqrt{(\cos\varphi)^2 + (\sin\varphi)^2 + (0)^2} = 1$   
•  $s_2 = \sqrt{\left(\frac{\partial x}{\partial \varphi}\right)^2 + \left(\frac{\partial y}{\partial \varphi}\right)^2 + \left(\frac{\partial z}{\partial \varphi}\right)^2} = \sqrt{(-\rho \sin\varphi)^2 + (\rho \cos\varphi)^2 + (0)^2} = \rho$   
•  $s_3 = \sqrt{\left(\frac{\partial x}{\partial z}\right)^2 + \left(\frac{\partial y}{\partial z}\right)^2 + \left(\frac{\partial z}{\partial z}\right)^2} = \sqrt{(0)^2 + (0)^2 + (1)^2} = 1$   
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8

#### Helmholtz Wave Equation भारतीय प्रौद्योगिकी संस्थान गुवाहाटी INDIAN INSTITUTE OF TECHNOLOGY GUWAHATI Ζ Spherical coordinate system θ • $a_1 = r, a_2 = \theta, a_3 = \varphi$ • $\hat{a}_1 = \hat{r}, \hat{a}_2 = \hat{\theta}, \hat{a}_3 = \hat{\varphi}$ V Relation between φ Cartesian and Spherical coordinates • $x = rsin\theta cos\phi$ , $y = rsin\theta sin\phi$ , $z = rcos\theta$ Calculation of scale factors • $s_1 = \sqrt{\left(\frac{\partial x}{\partial r}\right)^2 + \left(\frac{\partial y}{\partial r}\right)^2 + \left(\frac{\partial z}{\partial r}\right)^2} = \sqrt{(\sin\theta\cos\varphi)^2 + (\sin\theta\sin\varphi)^2 + (\cos\theta)^2} = 1$ • $s_2 = \sqrt{\left(\frac{\partial x}{\partial \theta}\right)^2 + \left(\frac{\partial y}{\partial \theta}\right)^2 + \left(\frac{\partial z}{\partial \theta}\right)^2} = \sqrt{(rcos\theta cos\phi)^2 + (rcos\theta sin\phi)^2 + (-rsin\theta)^2} = r$ • $s_3 = \sqrt{\left(\frac{\partial x}{\partial \varphi}\right)^2 + \left(\frac{\partial y}{\partial \varphi}\right)^2 + \left(\frac{\partial z}{\partial \varphi}\right)^2} = \sqrt{(-rsin\theta sin\varphi)^2 + (rsin\theta cos\varphi)^2 + (0)^2} = rsin\theta$

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• Laplacian of a scalar function:

$$\nabla^2 \psi = \nabla \bullet \nabla \psi = \frac{1}{s_1 s_2 s_3} \left[ \frac{\partial}{\partial a_1} \left( \frac{s_2 s_3}{s_1} \frac{\partial \psi}{\partial a_1} \right) + \frac{\partial}{\partial a_2} \left( \frac{s_1 s_3}{s_2} \frac{\partial \psi}{\partial a_2} \right) + \frac{\partial}{\partial a_3} \left( \frac{s_1 s_2}{s_3} \frac{\partial \psi}{\partial a_3} \right) \right]$$

#### Table: General Curvilinear System: Particular cases

General curvilinear system	a	(a <sub>2</sub> ) (	a <sub>3</sub>	(S <sub>1</sub> )	(S <sub>2</sub> )	S <sub>3</sub>
Cartesian coordinate system	X	(Y)	Ζ			
Cylindrical coordinate system	ρ	φ	(Z) (	1	ρ	
Spherical coordinate system (	$\overline{\mathbf{r}}$ (	θ	φ (	1	r	rsin

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$$\boldsymbol{\psi} = \nabla \bullet \nabla \boldsymbol{\psi}$$
$$= \boxed{\frac{1}{s_1 s_2 s_3}} \left[ \frac{\partial}{\partial a_1} \left( \underbrace{s_2 s_3}_{s_1} \frac{\partial \boldsymbol{\psi}}{\partial a_1} \right) + \frac{\partial}{\partial a_2} \left( \underbrace{s_1 s_3}_{s_2} \frac{\partial \boldsymbol{\psi}}{\partial a_2} \right) + \frac{\partial}{\partial a_3} \left( \underbrace{s_1 s_2}_{s_3} \frac{\partial \boldsymbol{\psi}}{\partial a_3} \right) \right]$$

#### Tips to memorize this formula:

- Note that in the expression of Laplacian of a scalar function above,
  - outside the third bracket, we have division by product of all scale factors, and
  - inside the third bracket there are three terms
- Each term is a partial differential with respect to a variable of the expression in a first bracket
- Inside first bracket,

 $\nabla^2 u$ 

- you have multiplication of scale factors of the remaining two axes divide by the scale factor of the same variable multiplied
  - partial differential of the scalar function with the same variable

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 $\nabla$ 

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 $\partial^2 \psi$ 

 $\partial z^2$ 

12

$$\begin{aligned} & \text{Laplacian of a scalar function:} \\ & \mathcal{P}^{2}\psi = \nabla \cdot \nabla \psi = \frac{1}{s_{1}s_{2}s_{3}} \left[ \frac{\partial}{\partial a_{1}} \left( \frac{s_{2}s_{3}}{s_{1}} \frac{\partial \psi}{\partial a_{1}} \right) + \frac{\partial}{\partial a_{2}} \left( \frac{s_{1}s_{3}}{s_{2}} \frac{\partial \psi}{\partial a_{2}} \right) + \frac{\partial}{\partial a_{3}} \left( \frac{s_{1}s_{2}}{s_{3}} \frac{\partial \psi}{\partial a_{3}} \right) \right] \\ & \text{Hence,} \\ & \text{Cartesian coordinates:} \quad \nabla^{2}\psi = \frac{\partial^{2}\psi}{\partial x^{2}} + \frac{\partial^{2}\psi}{\partial y^{2}} + \frac{\partial^{2}\psi}{\partial z^{2}} \\ & \text{Cylindrical coordinates:} \quad \nabla^{2}\psi = \frac{1}{\rho} \left[ \frac{\partial}{\partial \rho} \left( \rho \frac{\partial \psi}{\partial \rho} \right) + \frac{\partial}{\partial \phi} \left( \frac{1}{\rho} \frac{\partial \psi}{\partial \phi} \right) + \frac{\partial}{\partial z} \left( \rho \frac{\partial \psi}{\partial z} \right) \right] \\ & = \frac{1}{r^{2}} \sin \theta \left[ \frac{\partial}{\partial \rho} \left( r^{2} \sin \theta \frac{\partial \psi}{\partial r} \right) + \frac{\partial}{\partial \theta} \left( r \sin \theta \frac{\partial \psi}{\partial \theta} \right) + \frac{\partial}{\partial \phi} \left( \frac{r}{r \sin \theta} \frac{\partial \psi}{\partial \theta} \right) + \frac{\partial}{\partial \phi} \left( \frac{r}{r \sin \theta} \frac{\partial \psi}{\partial \theta} \right) \\ & = \frac{1}{r^{2}} \left[ \frac{\partial}{\partial r} \left( r^{2} \frac{\partial \psi}{\partial r} \right) \right] + \frac{1}{r^{2}} \sin \theta \left[ \sin \theta \frac{\partial \psi}{\partial \theta} \right] + \frac{1}{r^{2} \sin^{2} \theta} \left[ \frac{\partial^{2}\psi}{\partial \phi^{2}} \right] \end{aligned}$$