

FDTD: An Introduction

- Wave propagates more slowly on the medium on the right and
 - the pulse length is shorter
- *FDTD simulation of TM wave propagation due to multiple sources* (fDTD_2d_TM_multi_source.m)
- Just like ABC, we can have **Mur's Boundary Condition**
- Consider E_z component of a TM wave

$$\left[\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} - \frac{1}{v_p^2} \frac{\partial^2}{\partial t^2} \right] E_z = 0$$
$$\Rightarrow \left[\frac{\partial^2}{\partial x^2} - \left(\frac{1}{v_p^2} \frac{\partial^2}{\partial t^2} - \frac{\partial^2}{\partial y^2} \right) \right] E_z = 0$$

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- It can be expressed as

$$\left[\frac{\partial}{\partial x} + \sqrt{\frac{1}{v_p^2} \frac{\partial^2}{\partial t^2} - \frac{\partial^2}{\partial y^2}} \right] \left[\frac{\partial}{\partial x} - \sqrt{\frac{1}{v_p^2} \frac{\partial^2}{\partial t^2} - \frac{\partial^2}{\partial y^2}} \right] E_z = 0$$

$$\Rightarrow \Psi^+ \Psi^- E_z = 0; \Psi^+ = \frac{\partial}{\partial x} + \sqrt{\frac{1}{v_p^2} \frac{\partial^2}{\partial t^2} - \frac{\partial^2}{\partial y^2}}; \Psi^- = \frac{\partial}{\partial x} - \sqrt{\frac{1}{v_p^2} \frac{\partial^2}{\partial t^2} - \frac{\partial^2}{\partial y^2}}$$

- This operator can be expressed as

$$\Psi^+ = \frac{\partial}{\partial x} + \frac{1}{v_p} \frac{\partial}{\partial t} \sqrt{1 - S^2}; S \equiv \frac{v_p \frac{\partial}{\partial y}}{\frac{\partial}{\partial t}}$$

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- From Taylor's series expansion

$$\sqrt{1-S^2} = 1 - \frac{S^2}{2} - \frac{S^4}{8} - \frac{S^6}{16} - \dots$$

- Simplest approximation, first-order Mur Boundary Condition gives one way wave equation

$$\Psi^+ E_z = \frac{\partial E_z}{\partial x} + \frac{1}{v_p} \frac{\partial E_z}{\partial t} \cong 0$$

$$\Psi^- E_z = \frac{\partial E_z}{\partial x} - \frac{1}{v_p} \frac{\partial E_z}{\partial t} \cong 0$$

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- How do you discretize it?

$$\frac{\partial E_z}{\partial t} = v_p \frac{\partial E_z}{\partial x} \rightarrow \frac{E_z|_{3/2}^{n+1} - E_z|_{3/2}^n}{\Delta t} = v_p \frac{E_z|_2^{n+1/2} - E_z|_1^{n+1/2}}{\Delta x}$$

- Note that in 1-D TM, E_z is defined at integer time and space step, hence, we take time and space average for available values

$$\begin{aligned} & \frac{\left(E_z|_1^{n+1} + E_z|_2^{n+1}\right) - \left(E_z|_1^n + E_z|_2^n\right)}{2\Delta t} \\ &= v_p \frac{\left(E_z|_2^{n+1} + E_z|_2^n\right) - \left(E_z|_1^{n+1} + E_z|_1^n\right)}{2\Delta x} \end{aligned}$$

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- Simplifying

$$\begin{aligned} & \left[\left(E_z|_1^{n+1} + E_z|_2^{n+1} \right) - \left(E_z|_1^n + E_z|_2^n \right) \right] \\ &= \frac{\Delta t v_p}{\Delta x} \left[\left(E_z|_2^{n+1} + E_z|_2^n \right) - \left(E_z|_1^{n+1} + E_z|_1^n \right) \right] \end{aligned}$$

- Rearranging

$$\begin{aligned} \left(1 + \frac{v_p \Delta t}{\Delta x} \right) E_z|_1^{n+1} &= \left(1 + \frac{v_p \Delta t}{\Delta x} \right) E_z|_2^n - \left(1 - \frac{v_p \Delta t}{\Delta x} \right) E_z|_2^{n+1} + \left(1 - \frac{v_p \Delta t}{\Delta x} \right) E_z|_1^n \\ \Rightarrow E_z|_1^{n+1} &= E_z|_2^n + \left(\frac{v_p \Delta t - \Delta x}{v_p \Delta t + \Delta x} \right) \left(E_z|_2^{n+1} - E_z|_1^n \right) \end{aligned}$$

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- For CFL, it reduces to

$$E_z|_1^{n+1} = E_z|_2^n; \frac{v_p \Delta t}{\Delta x} = 1$$

- Similarly for the rightmost edge

$$E_z|_{i_{\max}}^{n+1} = E_z|_{i_{\max}-1}^n + \left(\frac{v_p \Delta t - \Delta x}{v_p \Delta t + \Delta x} \right) \left(E_z|_{i_{\max}-1}^{n+1} - E_z|_{i_{\max}}^n \right)$$

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Second-order Mur Boundary Condition keeps two terms of Taylor series expansion

$$\sqrt{1-S^2} \cong 1 - \frac{S^2}{2}$$

- Therefore,
$$\Psi^+ = \frac{\partial}{\partial x} + \frac{1}{v_p} \frac{\partial}{\partial t} \left(1 - \frac{S^2}{2} \right)$$

$$\Rightarrow \Psi^+ = \frac{\partial}{\partial x} + \frac{1}{v_p} \frac{\partial}{\partial t} - \frac{1}{2} \frac{1}{v_p} \frac{\partial}{\partial t} \left(\frac{v_p \frac{\partial}{\partial y}}{\frac{\partial}{\partial t}} \right)^2$$

$$\Rightarrow \Psi^+ = \frac{\partial}{\partial x} + \frac{1}{v_p} \frac{\partial}{\partial t} - \frac{v_p}{2} \left(\frac{\frac{\partial^2}{\partial y^2}}{\frac{\partial}{\partial t}} \right)$$

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- Hence,

$$\Psi^+ E_z = \frac{\partial E_z}{\partial x} + \frac{1}{v_p} \frac{\partial E_z}{\partial t} - \frac{v_p}{2} \left(\frac{\frac{\partial^2 E_z}{\partial y^2}}{\frac{\partial}{\partial t}} \right) = 0$$

$$\Rightarrow \frac{\partial^2 E_z}{\partial x \partial t} + \frac{1}{v_p} \frac{\partial^2 E_z}{\partial t^2} - \frac{v_p}{2} \frac{\partial^2 E_z}{\partial y^2} = 0$$

- Similarly,

$$\Psi^- E_z = 0$$

$$\Rightarrow \frac{\partial^2 E_z}{\partial x \partial t} - \frac{1}{v_p} \frac{\partial^2 E_z}{\partial t^2} + \frac{v_p}{2} \frac{\partial^2 E_z}{\partial y^2} = 0$$

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- How to discretize it?

$$\begin{aligned}
 \frac{\partial^2 E_z}{\partial x \partial t} \Big|_{3/2,j}^n &= \frac{1}{2\Delta t} \left(\frac{\partial E_z}{\partial x} \Big|_{3/2,j}^{n+1} - \frac{\partial E_z}{\partial x} \Big|_{3/2,j}^{n-1} \right) \\
 &= \frac{1}{2\Delta t} \left(\frac{E_z \Big|_{2,j}^{n+1} - E_z \Big|_{1,j}^{n+1}}{\Delta x} - \frac{E_z \Big|_{2,j}^{n-1} - E_z \Big|_{1,j}^{n-1}}{\Delta x} \right) \\
 \frac{\partial^2 E_z}{\partial y^2} \Big|_{3/2,j}^n &\cong \frac{1}{2} \left(\frac{\partial^2 E_z}{\partial y^2} \Big|_{2,j}^n + \frac{\partial^2 E_z}{\partial y^2} \Big|_{1,j}^n \right) \\
 &= \frac{1}{2} \left(\frac{E_z \Big|_{2,j+1}^n - 2E_z \Big|_{2,j}^n + E_z \Big|_{2,j-1}^n}{(\Delta y)^2} - \frac{E_z \Big|_{1,j+1}^n - 2E_z \Big|_{1,j}^n + E_z \Big|_{1,j-1}^n}{(\Delta y)^2} \right)
 \end{aligned}$$

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- Similarly

$$\begin{aligned}\frac{\partial^2 E_z}{\partial t^2}\bigg|_{3/2,j}^n &\cong \frac{1}{2} \left(\frac{\partial^2 E_z}{\partial t^2}\bigg|_{2,j}^n + \frac{\partial^2 E_z}{\partial t^2}\bigg|_{1,j}^n \right) \\ &= \frac{1}{2} \left(\frac{E_z|_{2,j}^{n+1} - 2E_z|_{2,j}^n + E_z|_{2,j}^{n-1}}{(\Delta t)^2} - \frac{E_z|_{1,j}^{n+1} - 2E_z|_{1,j}^n + E_z|_{1,j}^{n-1}}{(\Delta t)^2} \right)\end{aligned}$$

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- Finally

$$E_z|_{1,j}^{n+1} = -E_z|_{2,j}^{n-1} - \left(\frac{\Delta x - v_p \Delta t}{v_p \Delta t + \Delta x} \right) \left(E_z|_{2,j}^{n+1} + E_z|_{1,j}^{n-1} \right) + \left(\frac{2\Delta x}{v_p \Delta t + \Delta x} \right) \left(E_z|_{1,j}^n + E_z|_{2,j}^n \right) \\ + \left(\frac{\Delta x (v_p \Delta t)^2}{2(\Delta y)^2 (v_p \Delta t + \Delta x)} \right) \left(E_z|_{1,j+1}^n - 2E_z|_{1,j}^n + E_z|_{1,j-1}^n + E_z|_{2,j+1}^n - 2E_z|_{2,j}^n + E_z|_{2,j-1}^n \right)$$

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3. 3-D FDTD

- From two Maxwell's curl equations,
 - we need to consider all 6 equations for 3-D case

$$\frac{\partial \vec{E}}{\partial t} = \frac{1}{\epsilon_0 \epsilon_r} \nabla \times \vec{H}; \quad \frac{\partial \vec{H}}{\partial t} = -\frac{1}{\mu_0 \mu_r} \nabla \times \vec{E}$$

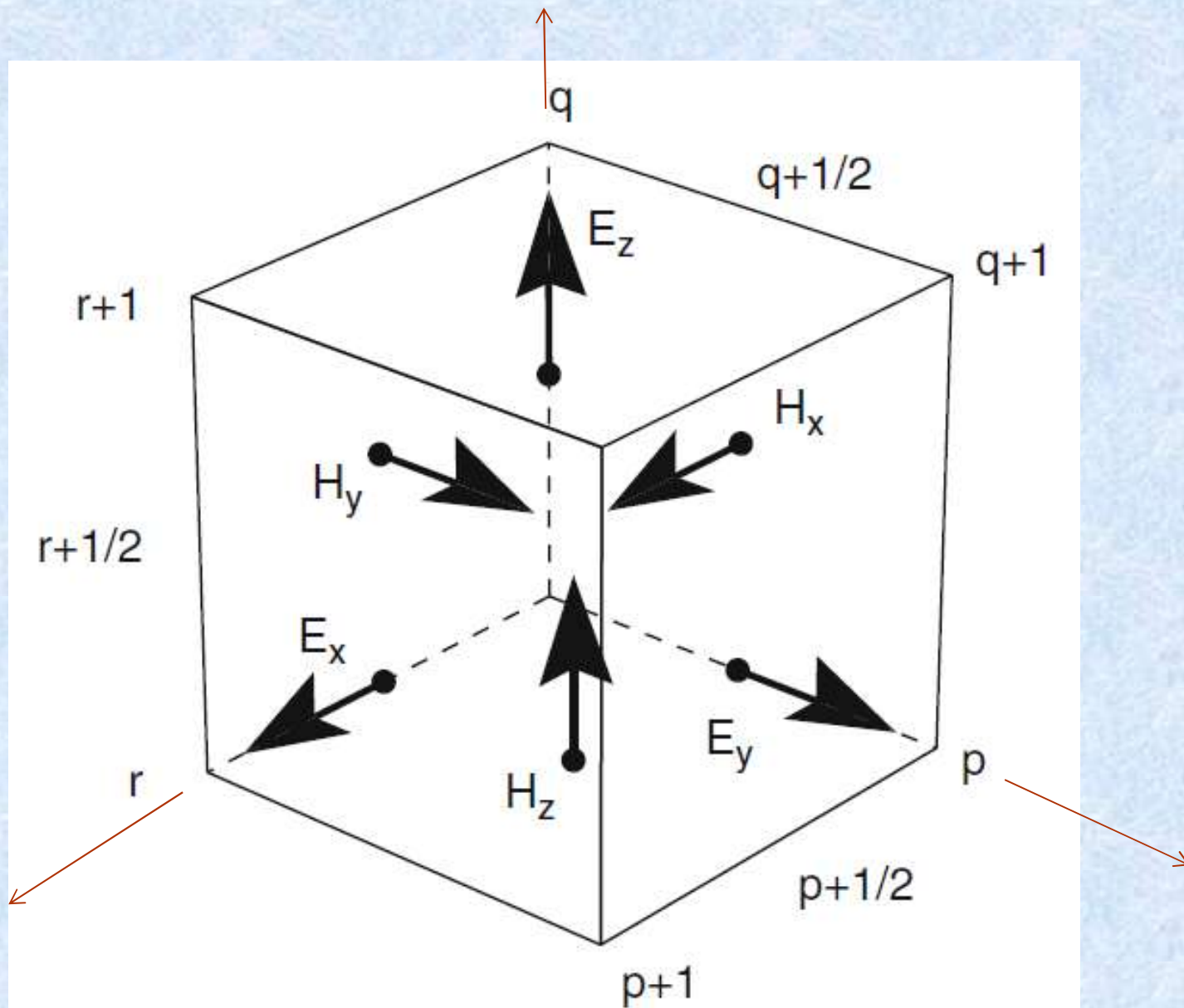


Fig. 4 Yee's FDTD grid

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Yee's FDTD grid

- Unlike 1-D and 2-D FDTD cases,
- Electric fields are calculated at “integer” time-steps
 - and magnetic field at “half-integer” time-steps
- Electric field components are placed at mid-points of the corresponding edges
- E.g.
 - E_x is placed at midpoints of edges oriented along x-direction
 - E_x is on half-grid in x and the integer grids in y & z

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- The magnetic field components are
 - placed at the centers of the faces of the cubes and
 - oriented normal to the faces
- E.g.
- H_x components are placed at the centers of the faces in the yz -plane
- Hence H_x is
 - on the integer grid in x and
 - on the half-grid in y & z

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Electric fields are calculated at
“integer” time-steps
and magnetic field at “half-integer” time-steps

E_x is placed at
midpoints of edges
oriented along x-
direction

$$\frac{\partial E_x}{\partial t} = \frac{1}{\epsilon_r \epsilon_0} \left(\frac{\partial H_z}{\partial y} - \frac{\partial H_y}{\partial z} \right)$$

H_z is placed at center
of faces in xy-plane

$$\frac{E_x^{n+1}(p+1/2, q, r) - E_x^n(p+1/2, q, r)}{\Delta t} = \frac{1}{\epsilon_r \epsilon_0} \left[\frac{H_z^{n+1/2}(p+1/2, q+1/2, r) - H_z^{n+1/2}(p+1/2, q-1/2, r)}{\Delta y} - \frac{H_y^{n+1/2}(p+1/2, q, r+1/2) - H_y^{n+1/2}(p+1/2, q, r-1/2)}{\Delta z} \right]$$

H_y is placed at center of faces in xz-plane

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$$\frac{\partial E_y}{\partial t} = \frac{1}{\epsilon_r \epsilon_0} \left(\frac{\partial H_x}{\partial z} - \frac{\partial H_z}{\partial x} \right)$$

$$\frac{E_y^{n+1}(p, q + 1/2, r) - E_y^n(p, q + 1/2, r)}{\Delta t} =$$
$$\frac{1}{\epsilon_r \epsilon_0} \left[\frac{H_x^{n+1/2}(p, q + 1/2, r + 1/2) - H_x^{n+1/2}(p, q + 1/2, r - 1/2)}{\Delta z} \right]$$
$$- \frac{1}{\epsilon_r \epsilon_0} \left[\frac{H_z^{n+1/2}(p + 1/2, q + 1/2, r) - H_z^{n+1/2}(p - 1/2, q + 1/2, r)}{\Delta z} \right]$$

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$$\frac{\partial E_z}{\partial t} = \frac{1}{\epsilon_r \epsilon_0} \left(\frac{\partial H_y}{\partial x} - \frac{\partial H_x}{\partial y} \right)$$

$$\begin{aligned} \frac{E_z^{n+1}(p, q, r+1/2) - E_z^n(p, q, r+1/2)}{\Delta t} = \\ \frac{1}{\epsilon_r \epsilon_0} \left[\frac{H_y^{n+1/2}(p+1/2, q, r+1/2) - H_y^{n+1/2}(p-1/2, q, r+1/2)}{\Delta x} \right] \\ - \frac{1}{\epsilon_r \epsilon_0} \left[\frac{H_x^{n+1/2}(p, q+1/2, r+1/2) - H_x^{n+1/2}(p, q-1/2, r+1/2)}{\Delta y} \right] \end{aligned}$$