

Introduction

- Euler's equation $\nabla^2 \Phi + k^2 \Phi = g$
- Functional is

$$I(\Phi) = \frac{1}{2} \iint \left[|\nabla \Phi|^2 - k^2 \Phi^2 + 2\Phi g \right] ds$$

- Euler's equation $\nabla^2 V = -\frac{\rho_v}{\varepsilon}$
- Transfer all terms to RHS $-\varepsilon \nabla^2 V - \rho_v = 0$
- Therefore, functional is

$$F(V_e) = \frac{1}{2} \int \left[\varepsilon |\nabla V_e|^2 - 2\rho_{V_e} V_e \right] ds$$

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- Energy functional for this case is

$$F(V_e) = \frac{1}{2} \int \left[\varepsilon |\nabla V_e|^2 - 2\rho_{V_e} V_e \right] ds$$

- which can be expressed as

$$F(V_e) = \frac{1}{2} \sum_{i=1}^3 \sum_{j=1}^3 \varepsilon V_{ei} \int \left[\nabla \alpha_i \cdot \nabla \alpha_j \right] ds V_{ej} \\ - \sum_{i=1}^3 \sum_{j=1}^3 V_{ei} \int \left[\alpha_i \alpha_j \right] ds \rho_{ej}$$

$$V_e = \sum_{i=1}^3 V_{ei} \alpha_i(x, y); \rho_{ve} = \sum_{i=1}^3 \rho_{ei} \alpha_i(x, y)$$

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$$F(V_e) = \frac{1}{2} \sum_{i=1}^3 \sum_{j=1}^3 \varepsilon V_{ei} \int [\nabla \alpha_i \bullet \nabla \alpha_j] ds V_{ej} - \sum_{i=1}^3 \sum_{j=1}^3 V_{ei} \int [\alpha_i \alpha_j] ds \rho_{ej}$$

- This can be written as

$$F(V_e) = \frac{1}{2} \varepsilon [V_e]^T [C^{(e)}] [V_e] - [V_e]^T [T^{(e)}] [\rho_e]$$

- where

$$C_{ij}^{(e)} = \int [\nabla \alpha_i \bullet \nabla \alpha_j] ds; T_{ij}^{(e)} = \int [\alpha_i \alpha_j] ds$$

- It can be shown that

$$T_{ij}^{(e)} = \begin{cases} A/12, & i \neq j \\ A/6, & i = j \end{cases}$$

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- Hence

$$F(V) = \sum_{e=1}^N F(V_e) = \frac{1}{2} \varepsilon [V]^T [C][V] - [V]^T [T][\rho]$$

- Solving the resulting equation:
- Iteration method:

$$\frac{\partial F(V)}{\partial V_k} = 0, k = 1, 2, \dots, n$$

- minimize the energy

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- For example for 5 node case

$$\frac{\partial F(V)}{\partial V_1} = \varepsilon \left[V_1 C_{11} + V_2 C_{21} + \dots + V_5 C_{51} \right]$$


$$- \left[T_{11} \rho_1 + \rho_2 T_{21} + \dots + T_{51} \rho_5 \right] = 0$$

- In general, for a free node k

$$V_k = -\frac{1}{C_{kk}} \sum_{i=1, j \neq k}^n V_i C_{ki} + \frac{1}{\varepsilon C_{kk}} \sum_{i=1}^n \rho_i T_{ki}$$

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- Band matrix method:

$$F(V) = \frac{1}{2} \varepsilon \begin{bmatrix} V_f & V_p \end{bmatrix} \begin{bmatrix} C_{ff} & C_{fp} \\ C_{pf} & C_{pp} \end{bmatrix} \begin{bmatrix} V_f \\ V_p \end{bmatrix} - \begin{bmatrix} V_f & V_p \end{bmatrix} \begin{bmatrix} T_{ff} & T_{fp} \\ T_{pf} & T_{pp} \end{bmatrix} \begin{bmatrix} \rho_f \\ \rho_p \end{bmatrix}$$



- Minimizing $F(V)$ w.r.t. V_f , we have,

$$\frac{\partial F}{\partial V_f} = 0$$

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- gives

$$\varepsilon \left(C_{ff} V_f + C_{pf} V_p \right) - \left(T_{ff} \rho_f + T_{fp} \rho_p \right) = 0$$

$$\Rightarrow [C_{ff}] [V_f] = - [C_{fp}] [V_p] + \frac{1}{\varepsilon} \left([T_{ff}] [\rho_f] + [T_{fp}] [\rho_p] \right)$$


- which can be written as


$$[A] [V] = [B]$$

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- where

$$[A] = [C_{ff}]$$

$$[V] = [V_f]$$

$$[B] = -[C_{fp}][V_p] + \frac{1}{\varepsilon} \left([T_{ff}][\rho_f] + [T_{fp}][\rho_p] \right)$$


- Solution of the wave equation:
- A typical wave equation is the inhomogeneous scalar Helmholtz equation

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- **Solution of inhomogeneous wave equation**

$$\nabla^2 \Phi + k^2 \Phi = g$$

- where Φ is the field quantity (for waveguide problem, $\Phi = H_z$ for TE mode or $\Phi = E_z$ for TM mode) to be determined, g is the source function and $k = \omega \sqrt{\mu \epsilon}$
- is the wavenumber of the medium
- (1) $k=0=g$; Laplace's equation
- (2) $k=0$; Poisson's equation
- (3) $g=0$, Helmholtz wave equation

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- Construction of functionals from PDEs
- For $L\Phi = g$ is obtained by extremizing the functional

$$I(\Phi) = \frac{1}{2} \iint \left[|\nabla\Phi|^2 - k^2\Phi^2 + 2\Phi g \right] ds$$

- Now express potential Φ and source function g in terms of shape functions α_i over a triangular element as

$$\Phi_e(x, y) = \sum_{i=1}^3 \alpha_i \Phi_{ei}$$

$$g_e(x, y) = \sum_{i=1}^3 \alpha_i g_{ei}$$

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- where Φ_{ei} and g_{ei} are respectively the values of Φ and g at nodal element i of element e
- Hence

$$I(\Phi_e) = \frac{1}{2} \sum_{i=1}^3 \sum_{j=1}^3 \Phi_{ei} \Phi_{ej} \int [\nabla \alpha_i \bullet \nabla \alpha_j] ds$$

$$- \frac{k^2}{2} \sum_{i=1}^3 \sum_{j=1}^3 \Phi_{ei} \Phi_{ej} \int [\alpha_i \alpha_j] ds + \sum_{i=1}^3 \sum_{j=1}^3 \Phi_{ei} g_{ej} \int [\alpha_i \alpha_j] ds$$

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$$I(\Phi_e) = \frac{1}{2}[\Phi_e]^T [C^{(e)}] [\Phi_e] - \frac{k^2}{2}[\Phi_e]^T [T^{(e)}] [\Phi_e] + [\Phi_e]^T [T^{(e)}] [G_e]$$

- where

$$[\Phi_e] = \begin{bmatrix} \Phi_{e1} \\ \Phi_{e2} \\ \Phi_{e3} \end{bmatrix}; [G_e] = \begin{bmatrix} g_{e1} \\ g_{e2} \\ g_{e3} \end{bmatrix}$$

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- Now

$$I(\Phi) = \sum_{e=1}^n I(\Phi_e) = \frac{1}{2}[\Phi]^T [C][\Phi] - \frac{k^2}{2}[\Phi]^T [T][\Phi] + [\Phi]^T [T][G]$$

- where

$$[\Phi] = \begin{bmatrix} \Phi_1 \\ \vdots \\ \Phi_N \end{bmatrix}; [G_e] = \begin{bmatrix} g_1 \\ \vdots \\ g_N \end{bmatrix}$$

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- [C] and [T] are global matrices consisting of local matrices
- [C^(e)] and [T^(e)] respectively
- When the source function $g=0$, if the free nodes are numbered first and prescribed nodes last, we have,

$$I(\Phi) = \frac{1}{2} \begin{bmatrix} \Phi_f & \Phi_p \end{bmatrix} \begin{bmatrix} C_{ff} & C_{fp} \\ C_{pf} & C_{pp} \end{bmatrix} \begin{bmatrix} \Phi_f \\ \Phi_p \end{bmatrix} - \frac{k^2}{2} \begin{bmatrix} \Phi_f & \Phi_p \end{bmatrix} \begin{bmatrix} T_{ff} & T_{fp} \\ T_{pf} & T_{pp} \end{bmatrix} \begin{bmatrix} \Phi_f \\ \Phi_p \end{bmatrix}$$

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- Setting $\frac{\partial I}{\partial \Phi_f} = 0$

$$\begin{bmatrix} C_{ff} & C_{fp} \end{bmatrix} \begin{bmatrix} \Phi_f \\ \Phi_p \end{bmatrix} - k^2 \begin{bmatrix} T_{ff} & T_{fp} \end{bmatrix} \begin{bmatrix} \Phi_f \\ \Phi_p \end{bmatrix} = 0$$

- For TM modes in a perfectly conducting waveguide, $\Phi_p = 0$

and hence

$$\begin{aligned} & \begin{bmatrix} C_{ff} & -k^2 T_{ff} \end{bmatrix} \Phi_f = 0 \\ & \Rightarrow \begin{bmatrix} T_{ff}^{-1} C_{ff} & -k^2 I \end{bmatrix} \Phi_f = 0 \\ & \Rightarrow (A - \lambda I)x = 0 \end{aligned}$$

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- where $A = T_{ff}^{-1} C_{ff}; \lambda = k^2; x = \Phi_f$
- It looks like an eigenvalue problem
- The eigenvalues are real since C and Y matrices are symmetric
- The above equation give eigenvalues and eigenfunctions of the Helmholtz problem