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Spectral domain MoM

Electric type dyadic Green's functions

$$\vec{E}(\vec{r}) = \int_V \vec{G}_{EJ}(\vec{r}, \vec{r}') \bullet \vec{J}_e(\vec{r}') d\vec{r}'$$

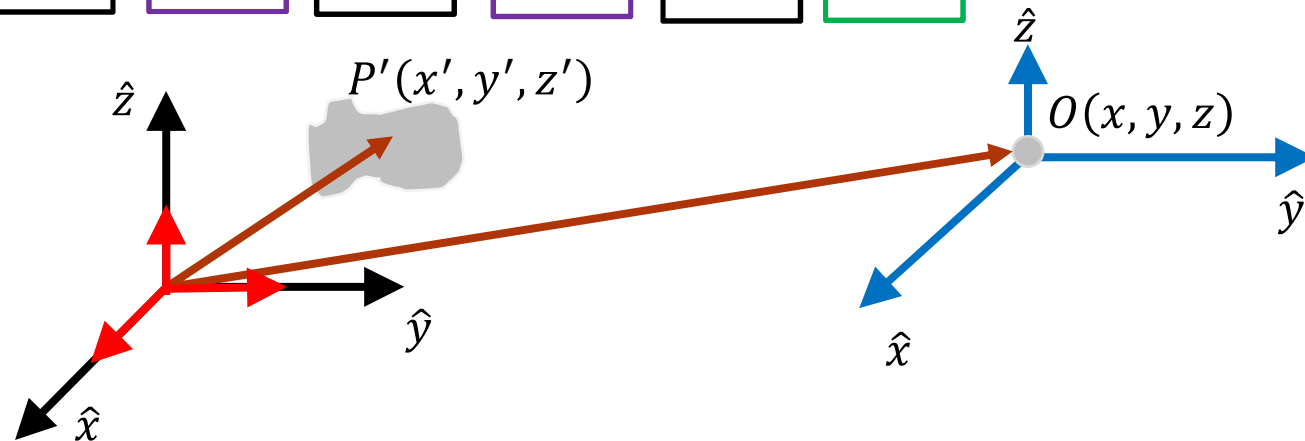
For 3-D current source, the electric dyadic Green's function will have nine components

$$\vec{G}_{EJ}(x, y, z; x', y', z') = G_{EJ}^{xx} \hat{x}\hat{x} + G_{EJ}^{xy} \hat{x}\hat{y} + G_{EJ}^{xz} \hat{x}\hat{z} + G_{EJ}^{yx} \hat{y}\hat{x} + G_{EJ}^{yy} \hat{y}\hat{y} + G_{EJ}^{yz} \hat{y}\hat{z} + G_{EJ}^{zx} \hat{z}\hat{x} + G_{EJ}^{zy} \hat{z}\hat{y} + G_{EJ}^{zz} \hat{z}\hat{z}$$

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- We are interested in first finding the spectral domain (Fourier transform domain) version
- of the above spatial domain dyadic Green's functions which will have nine components as follows

$$\tilde{\mathbf{G}}_{EJ}(\vec{k}_t, z, z') = \tilde{G}_{EJ}^{xx} \hat{x}\hat{x} + \tilde{G}_{EJ}^{xy} \hat{x}\hat{y} + \tilde{G}_{EJ}^{yx} \hat{y}\hat{x} + \tilde{G}_{EJ}^{yy} \hat{y}\hat{y} + \tilde{G}_{EJ}^{zx} \hat{z}\hat{x} + \tilde{G}_{EJ}^{xz} \hat{x}\hat{z} + \tilde{G}_{EJ}^{zy} \hat{z}\hat{y} + \tilde{G}_{EJ}^{yz} \hat{y}\hat{z} + \tilde{G}_{EJ}^{zz} \hat{z}\hat{z}$$



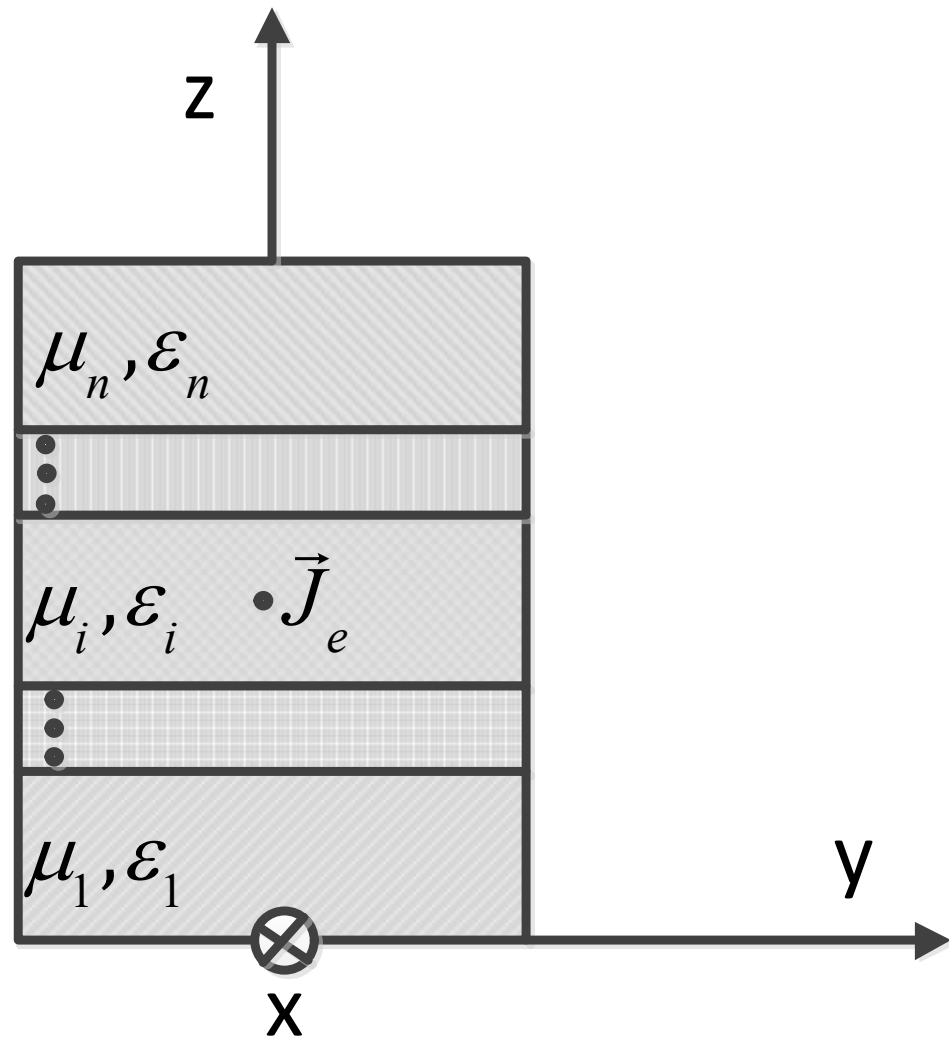
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- Usually, one neglects the metal thickness in the analysis which is applicable for perfect electric conductor
- For horizontal current excitation in x-y plane, removing the Green's functions which are dependent on the vertical current excitation component,
- only four Green's function components need to be determined

$$\tilde{\mathbf{G}}_{EJ}(\vec{k}_t, z, z') = \tilde{G}_{EJ}^{xx} \hat{x}\hat{x} + \tilde{G}_{EJ}^{xy} \hat{x}\hat{y} + \tilde{G}_{EJ}^{yx} \hat{y}\hat{x} + \tilde{G}_{EJ}^{yy} \hat{y}\hat{y}$$

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- Fig. One dimensional multilayer structures (n layer of dielectric and magnetic media) excited by three dimensional electric point source



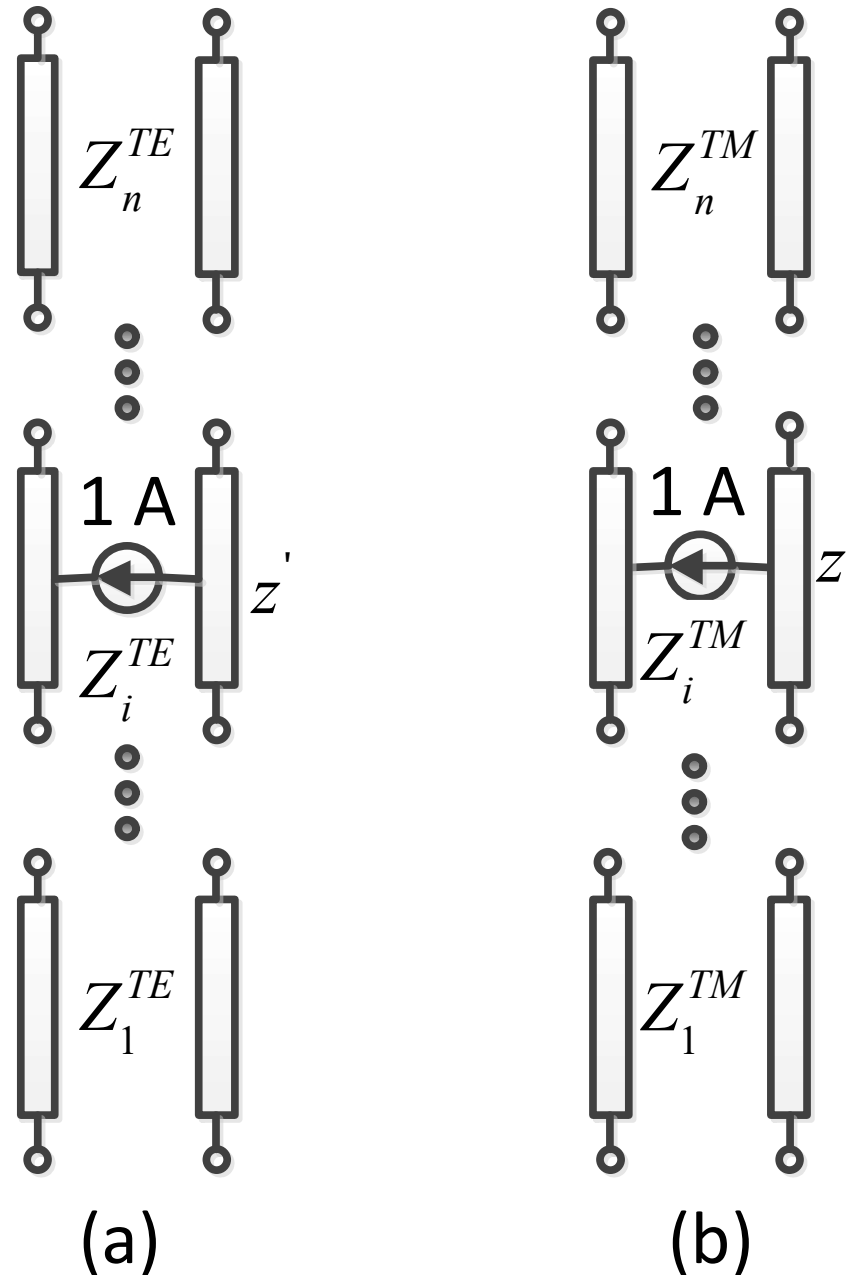
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- The four components of the spectral electric type transverse dyadic Green's function
- can be obtained from the voltages on the transverse electric (TE)/ transverse magnetic (TM) circuit models shown in Fig. as

$$\tilde{\tilde{\mathbf{G}}}_{EJ}(\vec{k}_t, z, z') = -V_{horizon}^{TE}(z, z')(\hat{k}_t \times \hat{z})(\hat{k}_t \times \hat{z}) - V_{horizon}^{TM}(z, z')(\hat{k}_t)(\hat{k}_t)$$

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- Fig. Equivalent TE/TM circuits models for one-dimensional multilayer structure ($i=1,2,3,\dots,n$) excited by a two-dimensional horizontal (transverse) electric point source (a) TE modes and (b) TM modes



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- using vector and dyadic analysis, we can simplify the above equation as follows

$$\begin{aligned}
 & \tilde{\mathbf{G}}_{EJ}(\vec{k}_t, z, z') \\
 &= -V_{horizon}^{TE}(z, z') \left(\frac{k_x \hat{x} + k_y \hat{y}}{\sqrt{k_x^2 + k_y^2}} \times \hat{z} \right) \left(\frac{k_x \hat{x} + k_y \hat{y}}{\sqrt{k_x^2 + k_y^2}} \times \hat{z} \right) - V_{horizon}^{TM}(z, z') \left(\frac{k_x \hat{x} + k_y \hat{y}}{\sqrt{k_x^2 + k_y^2}} \right) \left(\frac{k_x \hat{x} + k_y \hat{y}}{\sqrt{k_x^2 + k_y^2}} \right) \\
 &= -V_{horizon}^{TE}(z, z') \left(\frac{k_y \hat{x} - k_x \hat{y}}{\sqrt{k_x^2 + k_y^2}} \right) \left(\frac{k_y \hat{x} - k_x \hat{y}}{\sqrt{k_x^2 + k_y^2}} \right) - V_{horizon}^{TM}(z, z') \left(\frac{k_x^2 \hat{x}\hat{x} + k_x k_y \hat{x}\hat{y} + k_y k_x \hat{y}\hat{x} + k_y^2 \hat{y}\hat{y}}{k_x^2 + k_y^2} \right) \\
 &= -V_{horizon}^{TE}(z, z') \left(\frac{k_y^2 \hat{x}\hat{x} - k_y k_x (\hat{x}\hat{y} + \hat{y}\hat{x}) + k_x^2 \hat{y}\hat{y}}{k_x^2 + k_y^2} \right) - V_{horizon}^{TM}(z, z') \left(\frac{k_x^2 \hat{x}\hat{x} + k_x k_y (\hat{x}\hat{y} + \hat{y}\hat{x}) + k_y^2 \hat{y}\hat{y}}{k_x^2 + k_y^2} \right)
 \end{aligned}$$

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- Finally

$$\begin{aligned}
 & \tilde{\mathbf{G}}_{EJ}(\vec{k}_t, z, z') \\
 &= -\frac{V_{horizon}^{TE}(z, z')k_y^2 + V_{horizon}^{TM}(z, z')k_x^2}{k_x^2 + k_y^2} \hat{x}\hat{x} + \frac{k_x k_y (V_{horizon}^{TE}(z, z') - V_{horizon}^{TM}(z, z'))}{k_x^2 + k_y^2} (\hat{x}\hat{y} + \hat{y}\hat{x}) \\
 & \quad - \frac{V_{horizon}^{TE}(z, z')k_x^2 + V_{horizon}^{TM}(z, z')k_y^2}{k_x^2 + k_y^2} \hat{y}\hat{y}
 \end{aligned}$$

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- We have

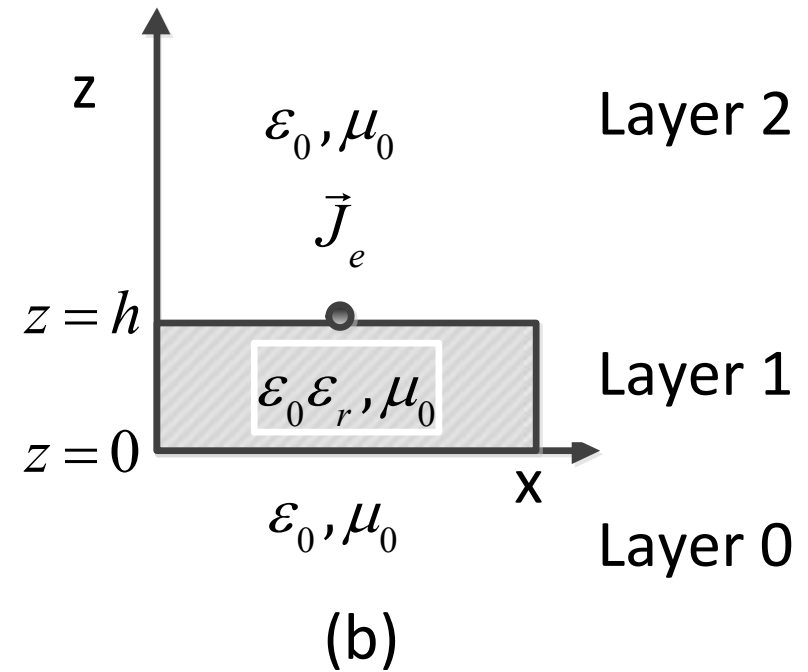
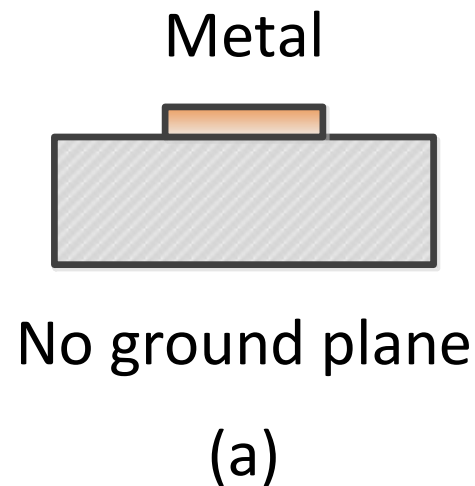
$$\tilde{G}_{EJ}^{xx}(k_x, k_y, z, z') = -\frac{V_{horizon}^{TE}(z, z')k_y^2 + V_{horizon}^{TM}(z, z')k_x^2}{k_x^2 + k_y^2}$$

$$\tilde{G}_{EJ}^{xy}(k_x, k_y, z, z') = \tilde{G}_{EJ}^{yx}(k_x, k_y, z, z') = \frac{k_x k_y (V_{horizon}^{TE}(z, z') - V_{horizon}^{TM}(z, z'))}{k_x^2 + k_y^2}$$

$$\tilde{G}_{EJ}^{yy}(k_x, k_y, z, z') = -\frac{V_{horizon}^{TE}(z, z')k_x^2 + V_{horizon}^{TM}(z, z')k_y^2}{k_x^2 + k_y^2}$$

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- Fig. (a) Side view of printed monopole antenna (PMA) (b) Transverse (horizontal) electric point source excitation in a three layered structure



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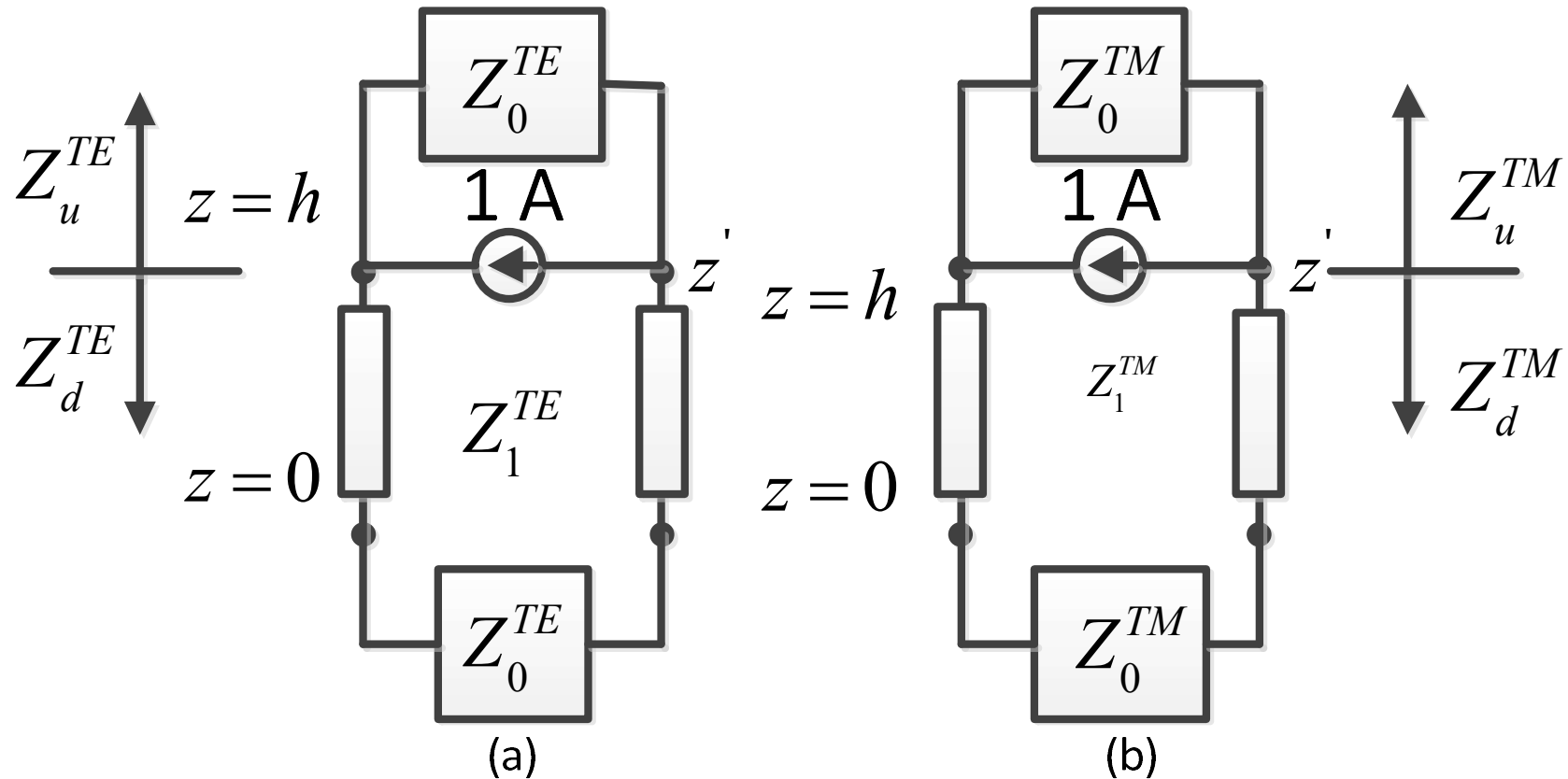


Fig. Equivalent (a) TE and (b) TM circuit models for horizontal (transverse) electric point source excitation

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- From the TE/TM circuit models, the impedance looking up the interface $z=h$ is simply the wave impedance of the free space

$$Z_u^{TE} = Z_0^{TE}; Z_u^{TM} = Z_0^{TM}$$

$$Z_0^{TE} = \frac{k_0 \eta_0}{\beta_0} \quad Z_0^{TM} = \frac{\beta_0 \eta_0}{k_0}$$

$$\beta_0 = \sqrt{k_0^2 - k_c^2} = \sqrt{k_0^2 - k_x^2 - k_y^2}$$

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- The impedance looking down the interface $z=h$ is simply given by a transmission line of length h terminated with the free space

$$Z_d^{TE} = Z_1^{TE} \frac{Z_0^{TE} + Z_1^{TE} \tanh(\gamma_1 h)}{Z_1^{TE} + Z_0^{TE} \tanh(\gamma_1 h)}; Z_d^{TM} = Z_1^{TM} \frac{Z_0^{TM} + Z_1^{TM} \tanh(\gamma_1 h)}{Z_1^{TM} + Z_0^{TM} \tanh(\gamma_1 h)}$$

$$Z_1^{TE} = \frac{k_1 \eta_1}{\beta_1} = \frac{k_0 \sqrt{\epsilon_r} \eta_0}{\beta_1 \sqrt{\epsilon_r}} = \frac{k_0 \eta_0}{\beta_1} \quad Z_1^{TM} = \frac{\beta_1 \eta_1}{k_1} = \frac{\beta_1 \eta_0}{k_0 \sqrt{\epsilon_r} \sqrt{\epsilon_r}} = \frac{\beta_1 \eta_0}{k_0 \epsilon_r}$$

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- In the above equation the equivalent impedance is the shunt combination of the impedance looking up and down at the interface
- Hence,

$$Z_{eq}^{TE} = \frac{Z_d^{TE} Z_u^{TE}}{Z_d^{TE} + Z_u^{TE}}; Z_{eq}^{TM} = \frac{Z_d^{TM} Z_u^{TM}}{Z_d^{TM} + Z_u^{TM}}$$

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- Finally, we may express the above expression simply

$$Z_{eq}^{TE} = k_0 \eta_0 \frac{1}{D_{TE}} \quad D_{TE} = \beta_0 + \frac{(\beta_0 + j\beta_1 \tan(\beta_1 h))}{\left(1 + j \frac{\beta_0}{\beta_1} \tan(\beta_1 h)\right)}$$

- Finally,

$$Z_{eq}^{TM} = \frac{\eta_0 \beta_0}{k_0 D_{TM}} \quad D_{TM} = 1 + \frac{\beta_0 \epsilon_r (\beta_1 + j\epsilon_r \beta_0 \tan(\beta_1 h))}{\beta_1 (\beta_0 \epsilon_r + j\beta_1 \tan(\beta_1 h))}$$

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- Green's function in terms of D_{TE} and D_{TM}

$$\tilde{G}_{EJ}^{xx}(\vec{k}_t, z, z' = h) = -\frac{1}{\omega \epsilon_0 k_\rho^2} \left(\frac{k_0^2 k_y^2}{D_{TE}} + \frac{\beta_0 k_x^2}{D_{TM}} \right)$$

$$\tilde{G}_{EJ}^{xy}(\vec{k}_t, z, z' = h) = \tilde{G}_{EJ}^{yx}(\vec{k}_t, z, z' = h) = \frac{k_x k_y}{\omega \epsilon_0 k_\rho^2} \left(\frac{k_0^2}{D_{TE}} - \frac{\beta_0}{D_{TM}} \right)$$

$$\tilde{G}_{EJ}^{yy}(\vec{k}_t, z, z' = h) = -\frac{1}{\omega \epsilon_0 k_\rho^2} \left(\frac{k_0^2 k_x^2}{D_{TE}} + \frac{\beta_0 k_y^2}{D_{TM}} \right)$$

$$k_\rho^2 = k_x^2 + k_y^2$$