Physics II (PH 102) Electromagnetism (Lecture 8 & 9)

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Electrostatic Potential Energy of a Point Test Charge in Electric Field

Consider a stationary configuration of source charge distribution.

- Let $\mathbf{E}(\mathbf{r})$ be a pre-existing electric field with potential $V(\mathbf{r})$ at a point P.
- ▶ The electrostatic force \mathbf{F}_{field} on a positive test charge Q is

 $\mathbf{F}_{field} = Q\mathbf{E} = -\mathbf{F}_{ext}$.

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Definition

Electrostatic Potential Energy of a test charge Q at point $P(\mathbf{r})$ is equal to the amount of work done by an external agent against the electrostatic field to bring the charge Q from ∞ (or ref) to the point $P(\mathbf{r})$:

$$U_E(\mathbf{r}) = \int_{-\infty}^{\mathbf{r}} \mathbf{F}_{ext} \cdot d\mathbf{r}' = -\int_{-\infty}^{\mathbf{r}} \mathbf{F}_{field} \cdot d\mathbf{r}' = -Q \int_{-\infty}^{\mathbf{r}} \mathbf{E} \cdot d\mathbf{r}' = QV(\mathbf{r}).$$

Note: The ambiguity in the absolute value of U_E at a point like V!

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Electrostatic Energy difference between two points a and b can be unambiguously expressed in terms of the Potential difference:

$$U_E(\mathbf{b}) - U_E(\mathbf{a}) = Q[V(\mathbf{b}) - V(\mathbf{a})] = -Q \int_{a}^{b} \mathbf{E}(\mathbf{r}') \cdot d\mathbf{r}'.$$

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Consider bringing in source charges $q_1, q_2, q_3, \dots, q_n$, one by one from ∞ .

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Consider bringing in source charges $q_1, q_2, q_3, \cdots, q_n$, one by one from ∞ .

- No work done in placing first charge q_1 at \mathbf{r}_1 , i.e., $W_1 = 0$
- Total Work done is placing up to the second charge q₂ at r₂:

$$W_2 = W_1 + \delta w_2 = 0 + \frac{1}{4\pi\epsilon_0} \frac{q_1q_2}{|\mathbf{r}_1 - \mathbf{r}_2|} = \frac{1}{4\pi\epsilon_0} \frac{q_1q_2}{r_{12}}$$



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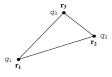
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Total Work done in placing up to the third charge q₃ at r₃:

$$W_3 = W_2 + \delta w_3 = W_2 + \frac{q_3}{4\pi\epsilon_0} \left(\frac{q_1}{r_{13}} + \frac{q_2}{r_{23}}\right)$$



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Consider bringing in source charges $q_1, q_2, q_3, \cdots, q_n$, one by one from ∞ .

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Total Work done in placing up to the third charge q₃ at r₃:

► Generalize formula up to *n* charges: $(j \neq i \text{ Self interactions excluded!})$ $W_n = W_{n-1} + \delta w_n = \frac{1}{4\pi\epsilon_0} \sum_{i=1}^{n-1} \sum_{j=2, j>i}^n \left(\frac{q_i q_j}{r_{ij}}\right) \rightarrow \text{No Double Counting}$ $\equiv \frac{1}{8\pi\epsilon_0} \sum_{i=1}^n \sum_{j=1, j\neq i}^n \left(\frac{q_i q_j}{r_{ij}}\right) \rightarrow \text{With Double Counting}$

▶ The Electrostatic Potential Energy is equal to the total work done W_n to assemble the configuration of n point charges at $\mathbf{r}_1, \mathbf{r}_2, \cdots, \mathbf{r}_n$:

$$U_E(\mathbf{r}_1, \mathbf{r}_2, \cdots, \mathbf{r}_n) \equiv W_n = \frac{1}{8\pi\epsilon_0} \sum_{i=1}^n \sum_{j=1, j\neq i}^n \left(\frac{q_i q_j}{r_{ij}}\right)$$
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$$= \frac{1}{2} \sum_{i=1}^n q_i V(\mathbf{r}_i) \equiv \frac{1}{2} \sum_{i=1}^n q_i V_i = \frac{1}{2} \sum_{i=1}^n u_{Ei}$$

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- ▶ $V_i \equiv V(\mathbf{r}_i)$: Potential at the *i*th charge site \mathbf{r}_i due to other n-1 charges.
- The result is independent of the order/sequence in which the charges are assembled at the respective EXACT locations, r₁, r₂, · · · , r_n.

$$U_{E}(\mathbf{r}_{1},\mathbf{r}_{2},\cdots,\mathbf{r}_{n})=U_{E}(\mathbf{r}_{2},\mathbf{r}_{n},\cdots,\mathbf{r}_{1})=\cdots=U_{E}\left[\operatorname{Permute}(\mathbf{r}_{1},\mathbf{r}_{2},\cdots,\mathbf{r}_{n})\right]$$

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The result OBVIOUSLY, depends on the respective locations of the charges. Thus, U_E is called the CONFIGURATION ENERGY.

Configuration Energy of Point Charges in free space

Superposition Principle is invalid (1/2 factor avoids double counting!):

$$U_E = \frac{1}{2} \sum_{i=1}^n u_{Ei}$$

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Recall:

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All terms i = j were absent in U_E .

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Recall:

- While defining U_E, we said "No work done in placing first charge q₁ at r₁."
- All terms i = j were absent in U_E .
- SELF-ENERGIES of the individual point charges were excluded!

Definition

SELF-ENERGY: This is the amount of energy needed to fabricate or build-up the individual point charges by bringing their respective differential amounts of constituent charges from ∞ to the specific locations.

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Electrostatic Potential Energy of General Distribution of Source Charges



(a) Discrete charges



(b) Line charge, λ



(c) Surface charge, σ



Localized charge distributions:

- Volume \mathcal{V} with volume charge density $ho(\mathbf{r})$
- Surface S with surface charge density $\sigma(\mathbf{r})$
- Curve Γ with linear charge density $\lambda(\mathbf{r})$
- Discrete point charges q_i at r'_i
- Potential V(r)
- Configuration Energy of system of charges:

 $U_{E} \approx \frac{1}{2} \iiint_{\mathcal{V}} \rho(\mathbf{r}') V(\mathbf{r}') d\tau' + \frac{1}{2} \iint_{S} \sigma(\mathbf{r}') V(\mathbf{r}') da'$ $+ \frac{1}{2} \int_{\Gamma} \lambda(\mathbf{r}') V(\mathbf{r}') dl' + \frac{1}{2} \sum_{i=1}^{n} q_{i} V(\mathbf{r}'_{i})$

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- Note: Self-energies of continuous distributions are included but not for discrete point charges!
- What else is missing here?

Electrostatic Potential Energy of General Distribution of Source Charges



(a) Discrete charges



(b) Line charge, λ



(c) Surface charge, σ



Localized charge distributions:

- Volume $\mathcal V$ with volume charge density $ho(\mathbf r)$
- Surface S with surface charge density $\sigma(\mathbf{r})$
- Curve Γ with linear charge density $\lambda(\mathbf{r})$
- Discrete point charges q_i at \mathbf{r}'_i
- Potential V(r)
- **Configuration Energy** of system of charges:

 $U_{E} \approx \frac{1}{2} \iiint_{\mathcal{V}} \rho(\mathbf{r}') V(\mathbf{r}') d\tau' + \frac{1}{2} \iint_{S} \sigma(\mathbf{r}') V(\mathbf{r}') da'$ $+ \frac{1}{2} \int_{\Gamma} \lambda(\mathbf{r}') V(\mathbf{r}') dl' + \frac{1}{2} \sum_{i=1}^{n} q_{i} V(\mathbf{r}'_{i})$

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- Note: Self-energies of continuous distributions are included but not for discrete point charges!
- What else is missing here? INTERACTION ENERGIES
- Superposition Principle is <u>invalid</u> in general!

True Configuration Energy of General Localized Charge Distribution

• Consider the most general localized charge distribution ρ_{tot} in a region \mathcal{V} , bounded by a closed surface S:

$$U_E = rac{1}{2} \iiint\limits_{\mathcal{V}}
ho_{ ext{tot}}(\mathbf{r}') V(\mathbf{r}') \, d au'.$$

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► Using Gauss's differential law <u>at source</u>: $\rho_{tot}(\mathbf{r}') = \epsilon_0 \nabla' \cdot \mathbf{E}(\mathbf{r}')$,

$$U_{E} = \frac{1}{2} \iiint_{\mathcal{V}} \rho(\mathbf{r}') V(\mathbf{r}') \, d\tau' = \frac{\epsilon_{0}}{2} \iiint_{\mathcal{V}} \left[\nabla' \cdot \mathbf{E}(\mathbf{r}') \right] V(\mathbf{r}') \, d\tau'$$

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$$U_{E} = \frac{1}{2} \iiint_{\mathcal{V}} \rho(\mathbf{r}') V(\mathbf{r}') d\tau' = \frac{\epsilon_{0}}{2} \iiint_{\mathcal{V}} [\nabla' \cdot \mathbf{E}(\mathbf{r}')] V(\mathbf{r}') d\tau'$$

$$= \frac{\epsilon_{0}}{2} \iiint_{\mathcal{V}} [|\mathbf{E}(\mathbf{r}')|^{2} + \nabla' \cdot [V(\mathbf{r}') \mathbf{E}(\mathbf{r}')]]_{\text{source}} d\tau'$$

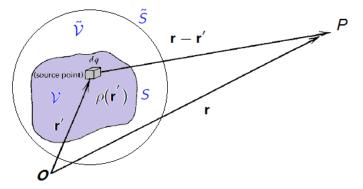
$$= \frac{\epsilon_{0}}{2} \iiint_{\mathcal{V}} |\mathbf{E}(\mathbf{r}')|^{2}_{\text{source}} d\tau' + \frac{\epsilon_{0}}{2} \oiint_{S} [V(\mathbf{r}') \mathbf{E}(\mathbf{r}')]_{\text{boundary}} \cdot d\mathbf{a}'.$$

In the last step, we applied the Gauss' Divergence Theorem to obtain the surface integral over the bounding surface S.

Recall Identity: $\nabla \cdot (V \mathbf{E}) = V (\nabla \cdot \mathbf{E}) + (\nabla V) \cdot \mathbf{E} = V (\nabla \cdot \mathbf{E}) - |\mathbf{E}|^2$

▶ Let us EXTEND the integration over a very large SPHERICAL volume $\tilde{\mathcal{V}}$, with bounding surface \tilde{S} , enclosing the LOCALIZED distribution \mathcal{V} :

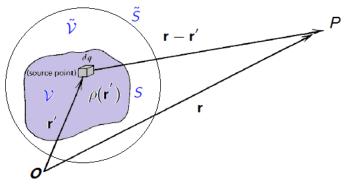
$$U_E = \frac{1}{2} \iiint_{\mathcal{V}} \rho_{\text{tot}}(\mathbf{r}') V(\mathbf{r}') \, d\tau' = \frac{1}{2} \iiint_{\hat{\mathcal{V}}} \rho_{\text{tot}}(\mathbf{r}') V(\mathbf{r}') \, d\tau'$$



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• For localized distribution at a distant point P, $V(\mathbf{r}) \propto 1/r$, $|\mathbf{E}(\mathbf{r})| \propto 1/r^2$ • For points \mathbf{r}' on this very large spherical surface \tilde{S} ,

$$\oint_{\tilde{S}} [V(\mathbf{r}') \mathbf{E}(\mathbf{r}')]_{\text{boundary}} \cdot d\mathbf{a}' \propto \left(\frac{1}{r'} \cdot \frac{1}{r'^2} \cdot r'^2\right)_{\tilde{S}} \sim \#\left(\frac{1}{r'}\right)_{\tilde{S}} \to 0, \ r' \to \infty.$$

True Configuration Energy of General Charge Distribution (contd.)

• Extending to include ALL SPACE: $\tilde{\mathcal{V}} \to \mathcal{V}_{\infty} \equiv \mathbb{R}^3$ and $\tilde{S} \to S_{\infty}$, then the surface integral vanishes!

$$U_{E} = \frac{1}{2} \iiint_{\mathcal{V}} \rho_{\text{tot}}(\mathbf{r}') V(\mathbf{r}') d\tau' = \frac{1}{2} \iiint_{\tilde{\mathcal{V}}} \rho_{\text{tot}}(\mathbf{r}') V(\mathbf{r}') d\tau'$$
$$= \frac{\epsilon_{0}}{2} \iiint_{\tilde{\mathcal{V}} \to \mathcal{V}_{\infty}} |\mathbf{E}|^{2}_{\text{source}} d\tau' + \frac{\epsilon_{0}}{2} \oiint_{\tilde{\mathcal{S}} \to \mathcal{S}_{\infty}} (V \mathbf{E})_{\text{boundary}} d\mathbf{a}'^{0}$$

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All inclusive formula: SELF-ENERGIES and INTERACTION ENERGIES of ALL localized charge distributions. True Configuration Energy of General Charge Distribution (contd.)

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$$= \frac{\epsilon_{0}}{2} \iiint_{\text{All Space}} |\mathbf{E}(\mathbf{r}')|^{2}_{\text{source}} d\tau'$$

All inclusive formula: SELF-ENERGIES and INTERACTION ENERGIES of ALL localized charge distributions.

In Summary: 2 Prescriptions to determine the Configuration Energy of a Localized distribution (depending on convenience) in a given problem

$$U_E = \frac{1}{2} \iiint_{\mathcal{V} \to \text{original vol.}} \rho_{\text{tot}}(\mathbf{r}) V(\mathbf{r}) \, d\tau = \frac{\epsilon_0}{2} \iiint_{\text{All Space}} |\mathbf{E}(\mathbf{r})|^2 \, d\tau$$

Configuration Energy of a Charged Sphere

Example

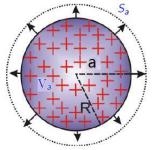
Determine the total configuration/self energy of a uniformly charged solid sphere V_R of radius R and charge q.

Configuration Energy of a Charged Sphere

Example

Determine the total configuration/self energy of a uniformly charged solid sphere V_R of radius R and charge q.

▶ AUGMENTED VOLUME: Consider a very large <u>concentric</u> spherical volume V_a of radius $a \gg R$, with its bounding surface S_a , enclosing the original charged sphere V_R .

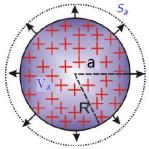


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$$U_{E} = \frac{1}{2} \iiint_{V_{R} \to V_{a}} \rho(\mathbf{r}) V(\mathbf{r}) d\tau = \frac{\epsilon_{0}}{2} \iiint_{V_{a}} |\mathbf{E}(\mathbf{r})|^{2} d\tau + \frac{\epsilon_{0}}{2} \oiint_{S_{a}} V(\mathbf{r}) \mathbf{E}(\mathbf{r}) \cdot d\mathbf{a}$$

Electric fields due to **original** charged sphere V_R (by Using Gauss's Law):

$$\mathsf{E}(\mathsf{r}) = \begin{cases} \frac{1}{4\pi\epsilon_0} \frac{qr}{R^3} \hat{\mathsf{r}} & r < R\\ \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \hat{\mathsf{r}} & r > R \end{cases}$$

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$$V(\mathbf{r}) = \begin{cases} \frac{q}{8\pi\epsilon_0 R} \left(3 - \frac{r^2}{R^2}\right) & r < R\\ \frac{1}{4\pi\epsilon_0} \frac{q}{r} & r > R \end{cases}$$

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▶ Volume integral over the **augmented** volume V_a of radius $a \gg R$:

$$\frac{\epsilon_0}{2} \iiint_{V_a} |\mathbf{E}(\mathbf{r})|^2 d\tau = \frac{\epsilon_0}{2} \left(\frac{q}{4\pi\epsilon_0}\right)^2 \left\{ \int_0^R \left(\frac{r^2}{R^6}\right) r^2 dr + \int_R^a \left(\frac{1}{r^4}\right) r^2 dr \right\} \int_0^{q_a} d\Omega$$
$$= \frac{q^2}{8\pi\epsilon_0} \left\{ \left(\frac{1}{5R}\right)_{r \le R} + \left(\frac{1}{R} - \frac{1}{a}\right)_{R < r \le a} \right\}$$

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Surface integral over the **augmented** sphere S_a of radius $a \gg R$:

$$\stackrel{\epsilon_0}{=} \oint_{S_a} \left[V(a\hat{\mathbf{r}}) \mathbf{E}(a\hat{\mathbf{r}}) \right] \cdot d\mathbf{a} = \frac{\epsilon_0}{2} \left(\frac{q}{4\pi\epsilon_0 a} \right) \left(\frac{q}{4\pi\epsilon_0 a^2} \right) \left(4\pi a^2 \right) = \left(\frac{q^2}{8\pi\epsilon_0 r} \right)_{r=a}$$

Extend V_a to include ALL SPACE:

$$V_a \to \mathcal{V}_\infty \equiv \mathbb{R}^3 \quad \& \quad S_a \to S_\infty \Longrightarrow a \to \infty$$

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$$= \frac{1}{4\pi\epsilon_{0}} \left(\frac{3q^{2}}{5R} \right)$$
$$U_{E} \equiv \frac{\epsilon_{0}}{2} \iiint_{\text{All Space}} |\mathbf{E}(\mathbf{r})|^{2} d\tau = \frac{1}{4\pi\epsilon_{0}} \left(\frac{3q^{2}}{5R} \right).$$

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Check by <u>direct</u> integration over the **original** charged sphere V_R:

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Electric field of a point charge placed at origin:

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► Self-Energy:

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$$= \frac{\epsilon_0}{2} \left(\frac{q}{4\pi\epsilon_0}\right)^2 \left\{ \int_0^\infty \left(\frac{1}{r^4}\right) 4\pi r^2 dr \right\}$$
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Since the radius δR of the point charge vanishes, the self-energy blow up!
 Point charges are <u>idealized</u> concepts. In reality δR ≠ 0 (say, for electrons), so Self-Energies (classically) of spherical objects of finite radius δR is:

$$W_{
m sphere} = rac{1}{4\pi\epsilon_0} \left(rac{3q^2}{5\,\delta R}
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Example

Find the Interaction energy of two charges, q_1 and q_2 located at r_1 and r_2 , respectively.

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 $\mathsf{E}(\mathsf{r})=\mathsf{E}_1(\mathsf{r})+\mathsf{E}_2(\mathsf{r})\,,$

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Total Configuration Energy:

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 \Rightarrow This is exactly the work done by an external agent in bringing q_2 from ∞ to r_2 with q_1 already present at r_1 .

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► It is conventional to define an **ENERGY DENSITY**:

$$u(\mathbf{r}) = \frac{\epsilon_0}{2} \left| \mathbf{E}(\mathbf{r}) \right|^2$$