Physics II (PH 102) Electromagnetism (Lecture 8 & 9)

Udit Raha

Indian Institute of Technology Guwahati

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Electrostatic Potential Energy of a Point Test Charge in Electric Field

Consider a stationary configuration of source charge distribution.

- In Let $E(r)$ be a pre-existing electric field with potential $V(r)$ at a point P.
- \blacktriangleright The electrostatic force \blacktriangleright \blacktriangleright and a positive test charge Q is

 $\mathbf{F}_{\text{field}} = \mathbf{OE} = -\mathbf{F}_{\text{ext}}$.

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 $F_{fold} = QE = -F_{ext}$.

Definition

Electrostatic Potential Energy of a test charge Q at point $P(r)$ is equal to the amount of work done by an external agent against the electrostatic field to bring the charge Q from ∞ (or ref) to the point $P(r)$:

$$
U_E(\mathbf{r}) = \int_{\infty}^{\mathbf{r}} \mathbf{F}_{ext} \cdot d\mathbf{r}' = -\int_{\infty}^{\mathbf{r}} \mathbf{F}_{field} \cdot d\mathbf{r}' = -Q \int_{\infty}^{\mathbf{r}} \mathbf{E} \cdot d\mathbf{r}' = QV(\mathbf{r}).
$$

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Note: The ambiguity in the absolute value of U_F at a point like V!

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Note: The ambiguity in the absolute value of U_F at a point like V!

 \blacktriangleright Electrostatic Energy difference between two points a and **b** can be unambiguously expressed in terms of the Potential difference:

$$
\overline{U_{E}(\mathbf{b})-U_{E}(\mathbf{a})}=Q[V(\mathbf{b})-V(\mathbf{a})]=-Q\int_{0}^{\mathbf{b}}\mathsf{E}(\mathbf{r}')\cdot d\mathbf{r}'.
$$

a

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Consider bringing in source charges $q_1, q_2, q_3, \cdots, q_n$, one by one from ∞ .

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Consider bringing in source charges $q_1, q_2, q_3, \cdots, q_n$, one by one from ∞ .

- \triangleright No work done in placing first charge q_1 at r_1 , i.e., $W_1 = 0$
- \blacktriangleright Total Work done is placing up to the second charge q_2 at r_2 :

$$
W_2 = W_1 + \delta w_2 = 0 + \frac{1}{4\pi\epsilon_0} \frac{q_1 q_2}{|r_1 - r_2|} = \frac{1}{4\pi\epsilon_0} \frac{q_1 q_2}{r_{12}}
$$

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 \blacktriangleright Total Work done in placing up to the third charge q_3 at r_3 :

$$
W_3 = W_2 + \delta w_3 = W_2 + \frac{q_3}{4\pi\epsilon_0} \left(\frac{q_1}{r_{13}} + \frac{q_2}{r_{23}}\right)
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*q*¹ **r¹** *q*² **r²** *q*³ **r³**

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$$

Generalize formula up to n charges: $(i \neq i$ Self interactions excluded!)

$$
W_n = W_{n-1} + \delta w_n = \frac{1}{4\pi\epsilon_0} \sum_{i=1}^{n-1} \sum_{j=2, j>i}^{n} \left(\frac{q_i q_j}{r_{ij}}\right) \to \text{No Double Counting}
$$

$$
\equiv \frac{1}{8\pi\epsilon_0} \sum_{i=1}^{n} \sum_{j=1, j\neq i}^{n} \left(\frac{q_i q_j}{r_{ij}}\right) \to \text{With Double Counting}
$$

 \blacktriangleright The Electrostatic Potential Energy is equal to the total work done W_n to assemble the configuration of *n* point charges at r_1, r_2, \cdots, r_n .

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U_E(\mathbf{r}_1, \mathbf{r}_2, \cdots, \mathbf{r}_n) \equiv W_n = \frac{1}{8\pi\epsilon_0} \sum_{i=1}^n \sum_{j=1, j\neq i}^n \left(\frac{q_i q_j}{r_{ij}}\right)
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=
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= $\frac{1}{2} \sum_{i=1}^{n} q_{i} V(\mathbf{r}_{i}) \equiv \frac{1}{2} \sum_{i=1}^{n} q_{i} V_{i} = \frac{1}{2} \sum_{i=1}^{n} u_{Ei}$

 $V_i \equiv V(r_i)$: Potential at the *ith* charge site r_i due to other $n-1$ charges.

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- $V_i \equiv V(r_i)$: Potential at the *ith* charge site r_i due to other $n-1$ charges.
- ▶ The result is independent of the order/sequence in which the charges are assembled at the respective EXACT locations, $\mathbf{r}_1, \mathbf{r}_2, \cdots, \mathbf{r}_n$.

$$
U_E(\mathbf{r}_1, \mathbf{r}_2, \cdots, \mathbf{r}_n) = U_E(\mathbf{r}_2, \mathbf{r}_n, \cdots, \mathbf{r}_1) = \cdots = U_E\left[\mathrm{Permute}(\mathbf{r}_1, \mathbf{r}_2, \cdots, \mathbf{r}_n)\right]
$$

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 U_E ($\mathbf{r}_1, \mathbf{r}_2, \cdots, \mathbf{r}_n$) = U_E ($\mathbf{r}_2, \mathbf{r}_n, \cdots, \mathbf{r}_1$) = $\cdots = U_E$ [Permute($\mathbf{r}_1, \mathbf{r}_2, \cdots, \mathbf{r}_n$)]

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 \triangleright The result OBVIOUSLY, depends on the respective locations of the charges. Thus, U_F is called the CONFIGURATION ENERGY.

Configuration Energy of Point Charges in free space

 \triangleright Superposition Principle is invalid (1/2 factor avoids double counting!):

$$
U_E = \frac{1}{2} \sum_{i=1}^n u_{E_i}
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Recall:

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All terms $i = j$ were absent in U_F .

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Recall:

- \triangleright While defining U_F , we said "No work done in placing first charge q_1 at r_1 ."
- All terms $i = j$ were absent in U_F .
- \triangleright SELF-ENERGIES of the individual point charges were excluded!

Definition

✎ SELF-ENERGY: This is the amount of energy needed to fabricate or build-up ✍ the individual point charges by bringing their respective differential amounts of constituent charges from ∞ to the specific locations.

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Electrostatic Potential Energy of General Distribution of Source Charges

(a) Discrete charges

(b) Line charge, λ

(c) Surface charge, σ

Localized charge distributions:

- \triangleright Volume $\mathcal V$ with volume charge density $\rho(\mathbf r)$
- Surface S with surface charge density $\sigma(\mathbf{r})$
- **I** Curve Γ with linear charge density $\lambda(r)$
- Discrete point charges q_i at r'_i
- \blacktriangleright Potential $V(r)$
- \triangleright Configuration Energy of system of charges:

 $U_E \approx \frac{1}{2}$ 2 $\iiint \rho(\mathbf{r}')V(\mathbf{r}') d\tau' + \frac{1}{2}$ $\mathcal V$ 2 $\iint \sigma(\mathbf{r}')V(\mathbf{r}') d\mathbf{a}'$ S $+\frac{1}{2}$ 2 ˆ Γ λ (r') $V(r')$ dl' + $\frac{1}{2}$ 2 $\sum_{n=1}^{\infty}$ $i=1$ $q_i V(\mathbf{r}'_i)$

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- \triangleright Note: Self-energies of continuous distributions are included but not for discrete point charges!
- \blacktriangleright What else is missing here?

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- \triangleright Note: Self-energies of continuous distributions are included but not for discrete point charges!
- \blacktriangleright What else is missing here? INTERACTION ENERGIES
- \triangleright Superposition Principle is invalid in general!

True Configuration Energy of General Localized Charge Distribution

In Consider the most general localized charge distribution ρ_{tot} **in a region** V **,** bounded by a closed surface S :

$$
U_E = \frac{1}{2} \iiint\limits_{\mathcal{V}} \rho_{\text{tot}}(\mathbf{r}') V(\mathbf{r}') d\tau'.
$$

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► Using Gauss's differential law <u>at source</u>: $\rho_{\rm tot}(\mathbf{r}') = \epsilon_0 \nabla' \cdot \mathbf{E}(\mathbf{r}'),$

$$
U_E = \frac{1}{2} \iiint\limits_{V} \rho(\mathbf{r}') V(\mathbf{r}') d\tau' = \frac{\epsilon_0}{2} \iiint\limits_{V} \left[\nabla' \cdot \mathbf{E}(\mathbf{r}') \right] V(\mathbf{r}') d\tau'
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$$

\n
$$
= \frac{\epsilon_{0}}{2} \iiint\limits_{V} \left[\left| \mathbf{E}(\mathbf{r}') \right|^{2} + \nabla' \cdot \left[V(\mathbf{r}') \mathbf{E}(\mathbf{r}') \right] \right]_{\text{source}} d\tau'
$$

\n
$$
= \frac{\epsilon_{0}}{2} \iiint\limits_{V} \left| \mathbf{E}(\mathbf{r}') \right|_{\text{source}}^{2} d\tau' + \frac{\epsilon_{0}}{2} \oiint\limits_{S} \left[V(\mathbf{r}') \mathbf{E}(\mathbf{r}') \right]_{\text{boundary}} \cdot d\mathbf{a}'.
$$

In the last step, we applied the Gauss' Divergence Theorem to obtain the surface integral over the bounding surface S .

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Recall Identity: $\nabla \cdot (V \mathbf{E}) = V(\nabla \cdot \mathbf{E}) + (\nabla V) \cdot \mathbf{E} = V(\nabla \cdot \mathbf{E}) - |\mathbf{E}|^2$

 \blacktriangleright Let us EXTEND the integration over a very large SPHERICAL volume $\tilde{\mathcal{V}}$, with bounding surface \tilde{S} , enclosing the LOCALIZED distribution \mathcal{V} :

$$
U_E = \frac{1}{2} \iiint\limits_{\mathcal{V}} \rho_{\text{tot}}(\mathbf{r}') V(\mathbf{r}') d\tau' = \frac{1}{2} \iiint\limits_{\tilde{\mathcal{V}}} \rho_{\text{tot}}(\mathbf{r}') V(\mathbf{r}') d\tau'
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$$

For localized distribution at a distant point P, $V(r) \propto 1/r$, $|E(r)| \propto 1/r^2$

 \blacktriangleright For points r' on this very large spherical surface \tilde{S}_i

$$
\oiint_{\tilde{S}} [V(\mathbf{r}') \mathbf{E}(\mathbf{r}')]_{\text{boundary}} \cdot d\mathbf{a}' \propto \left(\frac{1}{r'} \cdot \frac{1}{r'^2} \cdot r'^2\right)_{\tilde{S}} \sim # \left(\frac{1}{r'}\right)_{\tilde{S}} \rightarrow 0, r' \rightarrow \infty.
$$

True Configuration Energy of General Charge Distribution (contd.)

▶ Extending to include <code>ALL SPACE</code>: $\tilde{\mathcal{V}} \to \mathcal{V}_\infty \equiv \mathbb{R}^3$ and $\tilde{S} \to S_\infty$, then the surface integral vanishes!

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$$
= \frac{\epsilon_0}{2} \iiint\limits_{\tilde{V} \to V_{\infty}} |\mathbf{E}|_{\text{source}}^2 d\tau' + \frac{\epsilon_0}{2} \oiint\limits_{\tilde{S} \to S_{\infty}} (V\mathbf{E})_{\text{boundary}} d\mathbf{a}'
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$$

\n
$$
= \frac{\epsilon_{0}}{2} \iiint\limits_{\text{All Space}} |\mathbf{E}(\mathbf{r}')|_{\text{source}}^{2} d\tau'
$$

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▶ All inclusive formula: SELF-ENERGIES and INTERACTION ENERGIES of ALL localized charge distributions.

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In Summary: 2 Prescriptions to determine the Configuration Energy of a Localized distribution (depending on convenience) in a given problem

$$
U_E=\frac{1}{2}\mathop{\iiint}\limits_{\mathcal{V}\rightarrow \text{original vol.}}\rho_{\text{tot}}(\textbf{r})\,V(\textbf{r})\,d\tau=\frac{\epsilon_0}{2}\mathop{\iiint}\limits_{\text{All Space}}|\textbf{E}(\textbf{r})|^2\,\,d\tau
$$

Configuration Energy of a Charged Sphere

Example

Determine the total configuration/self energy of a uniformly charged solid sphere V_R of radius R and charge q .

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AUGMENTED VOLUME: Consider a very large concentric spherical volume V_a of radius $a \gg R$, with its bounding surface S_a , enclosing the original charged sphere V_R

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Determine the total configuration/self energy of a uniformly charged solid sphere V_R of radius R and charge q .

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$$
U_E = \frac{1}{2} \iiint\limits_{V_R \to V_a} \rho(\mathbf{r}) V(\mathbf{r}) d\tau = \frac{\epsilon_0}{2} \iiint\limits_{V_a} |\mathbf{E}(\mathbf{r})|^2 d\tau + \frac{\epsilon_0}{2} \oiint\limits_{S_a} V(\mathbf{r}) \mathbf{E}(\mathbf{r}) \cdot d\mathbf{a}
$$

$$
\mathsf{E}(r) = \begin{cases} \frac{1}{4\pi\epsilon_0} \frac{qr}{R^3} \hat{r} & r < R\\ \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \hat{r} & r > R \end{cases}
$$

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\mathsf{E}(\mathsf{r}) = \begin{cases} \frac{1}{4\pi\epsilon_0} \frac{q\mathsf{r}}{R^3} \hat{\mathsf{r}} & \mathsf{r} < R\\ \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \hat{\mathsf{r}} & \mathsf{r} > R \end{cases}
$$

▶ Potentials due to **original** charged sphere V_R (using $V(r) = -\int_{\infty}^{r} \mathbf{E} \cdot d\mathbf{l}$):

$$
V(r) = \begin{cases} \frac{q}{8\pi\epsilon_0 R} \left(3 - \frac{r^2}{R^2}\right) & r < R\\ \frac{1}{4\pi\epsilon_0} \frac{q}{r} & r > R \end{cases}
$$

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$$
\mathsf{E}(\mathsf{r}) = \begin{cases} \frac{1}{4\pi\epsilon_0} \frac{q\mathsf{r}}{R^3} \hat{\mathsf{r}} & \mathsf{r} < R\\ \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \hat{\mathsf{r}} & \mathsf{r} > R \end{cases}
$$

▶ Potentials due to **original** charged sphere V_R (using $V(r) = -\int_{\infty}^{r} \mathbf{E} \cdot d\mathbf{l}$):

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$$

 \blacktriangleright Volume integral over the augmented volume V_a of radius $a \gg R$:

$$
\frac{\epsilon_0}{2} \iiint\limits_{V_a} |\mathbf{E}(\mathbf{r})|^2 d\tau = \frac{\epsilon_0}{2} \left(\frac{q}{4\pi\epsilon_0} \right)^2 \left\{ \int_0^R \left(\frac{r^2}{R^6} \right) r^2 dr + \int_R^a \left(\frac{1}{r^4} \right) r^2 dr \right\} \int_0^{4\pi} d\Omega
$$

$$
= \frac{q^2}{8\pi\epsilon_0} \left\{ \left(\frac{1}{5R} \right)_{r \leq R} + \left(\frac{1}{R} - \frac{1}{a} \right)_{R < r \leq a} \right\}
$$

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$$
\mathsf{E}(\mathsf{r}) = \begin{cases} \frac{1}{4\pi\epsilon_0} \frac{q\mathsf{r}}{R^3} \hat{\mathsf{r}} & \mathsf{r} < R\\ \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \hat{\mathsf{r}} & \mathsf{r} > R \end{cases}
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$$

 \blacktriangleright Volume integral over the augmented volume V_a of radius $a \gg R$:

$$
\frac{\epsilon_0}{2} \iiint\limits_{V_3} |\mathbf{E}(\mathbf{r})|^2 d\tau = \frac{\epsilon_0}{2} \left(\frac{q}{4\pi\epsilon_0} \right)^2 \left\{ \int_0^R \left(\frac{r^2}{R^6} \right) r^2 dr + \int_R^a \left(\frac{1}{r^4} \right) r^2 dr \right\} \int_0^{4\pi} d\Omega
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$$

Surface integral over the **augmented** sphere S_a of radius $a \gg R$:

$$
\frac{\epsilon_0}{2} \oiint\limits_{S_a} [V(a\hat{\mathbf{r}}) \mathbf{E}(a\hat{\mathbf{r}})] \cdot d\mathbf{a} = \frac{\epsilon_0}{2} \left(\frac{q}{4\pi\epsilon_0 a} \right) \left(\frac{q}{4\pi\epsilon_0 a^2} \right) (4\pi a^2) = \left(\frac{q^2}{8\pi\epsilon_0 r} \right)_{r=a}
$$

$$
V_a \to \mathcal{V}_{\infty} \equiv \mathbb{R}^3 \quad \& \quad S_a \to S_{\infty} \Longrightarrow a \to \infty
$$

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D Configuration energy (self-energy) of the charged sphere V_R :

$$
U_E = \frac{\epsilon_0}{2} \iiint\limits_{V_a \to V_{\infty}} |E(\mathbf{r})|^2 d\tau + \frac{\epsilon_0}{2} \iiint\limits_{S_a \to S_{\infty}} V(\mathbf{r}) E(\mathbf{r}) \cdot d\mathbf{a}
$$

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U_{E} = \frac{\epsilon_{0}}{2} \iiint\limits_{V_{a} \to V_{\infty}} |E(r)|^{2} d\tau + \frac{\epsilon_{0}}{2} \iiint\limits_{S_{a} \to S_{\infty}} V(r) E(r) \cdot d\mathbf{a}
$$

\n
$$
= \lim_{a \to \infty} \left[\frac{q^{2}}{8\pi \epsilon_{0}} \left\{ \frac{1}{5R} + \left(\frac{1}{R} - \frac{1}{A} \right) \right\} + \frac{q^{2}}{8\pi \epsilon_{0} a} \right]^{0}
$$

\n
$$
= \frac{1}{4\pi \epsilon_{0}} \left(\frac{3q^{2}}{5R} \right)
$$

\n
$$
U_{E} \equiv \frac{\epsilon_{0}}{2} \iiint\limits_{\text{All Space}} |E(r)|^{2} d\tau = \frac{1}{4\pi \epsilon_{0}} \left(\frac{3q^{2}}{5R} \right).
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$$

 \blacktriangleright Check by direct integration over the original charged sphere V_R :

$$
U_E = \frac{1}{2} \iiint\limits_{V_R} \rho(\mathbf{r}) V(\mathbf{r}) d\tau = \frac{\epsilon_0}{2} \left(\frac{q}{4\pi\epsilon_0} \right)^2 \int\limits_0^R \left(\frac{r^2}{R^6} \right) r^2 dr \int\limits_0^{4\pi} d\Omega = \frac{1}{4\pi\epsilon_0} \left(\frac{3q^2}{5R} \right).
$$

 \blacktriangleright Electric field of a point charge placed at origin:

$$
\mathsf{E}(\mathsf{r})=\frac{1}{4\pi\epsilon_0}\frac{q}{r^2}\hat{\mathsf{r}}
$$

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 \blacktriangleright Self-Energy:

 \blacktriangleright Electric field of a point charge placed at origin:

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$$
W = \frac{\epsilon_0}{2} \iiint\limits_{\text{All Space}} |\mathbf{E}(\mathbf{r})|^2 d\tau = \frac{\epsilon_0}{2} \left(\frac{q}{4\pi\epsilon_0}\right)^2 \left\{ \int_0^\infty \left(\frac{1}{r^4}\right) r^2 dr \int_0^\pi \sin\theta d\theta \int_0^{2\pi} d\phi \right\}
$$

$$
= \frac{\epsilon_0}{2} \left(\frac{q}{4\pi\epsilon_0}\right)^2 \left\{ \int_0^\infty \left(\frac{1}{r^4}\right) 4\pi r^2 dr \right\}
$$

$$
= \frac{q^2}{8\pi\epsilon_0} \int_0^\infty \frac{1}{r^2} dr
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Since the radius δR of the point charge vanishes, the self-energy blow up!

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= \frac{\epsilon_0}{2} \iiint\limits_{\text{Bpace}} |\mathbf{E}(\mathbf{r})|^2 d\tau = \frac{\epsilon_0}{2} \left(\frac{q}{4\pi\epsilon_0}\right)^2 \left\{ \int_0^\infty \left(\frac{1}{r^4}\right) r^2 dr \int_0^{\pi} \sin\theta d\theta \int_0^{2\pi} d\phi \right\}
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$$

Since the radius δR of the point charge vanishes, the self-energy blow up! Point charges are idealized concepts. In reality $\delta R \neq 0$ (say, for electrons), so Self-Energies (classically) of spherical objects of finite radius δR is:

$$
W_{\rm sphere} = \frac{1}{4\pi\epsilon_0} \left(\frac{3q^2}{5 \,\delta R}\right) \rightarrow \text{finite}
$$

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Example

Find the Interaction energy of two charges, q_1 and q_2 located at r_1 and r_2 , respectively.

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Example

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Net Electric field at any point r (Superposition Principle)

 $E(r) = E_1(r) + E_2(r),$

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Total Configuration Energy:

$$
W = \frac{\epsilon_0}{2} \iiint\limits_{\text{All Space}} |\mathbf{E}(\mathbf{r})|^2 d\tau = \frac{\epsilon_0}{2} \iiint\limits_{\text{All Space}} |\mathbf{E}_1(\mathbf{r}) + \mathbf{E}_2(\mathbf{r})|^2 d\tau
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$$

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= \frac{\epsilon_0}{2} \iiint\limits_{\text{All Space}} |\mathbf{E}_1(\mathbf{r})|^2 d\tau + \frac{\epsilon_0}{2} \iiint\limits_{\text{All Space}} |\mathbf{E}_2(\mathbf{r})|^2 d\tau + \frac{\epsilon_0}{2} \iiint\limits_{\text{All Space}} 2\mathbf{E}_1(\mathbf{r}) \cdot \mathbf{E}_2(\mathbf{r}) d\tau
$$

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$$

=
$$
\frac{\epsilon_0}{2} \iiint\limits_{\text{All Space}} |E_1(\mathbf{r})|^2 d\tau + \frac{\epsilon_0}{2} \iiint\limits_{\text{All Space}} |E_2(\mathbf{r})|^2 d\tau + \frac{\epsilon_0}{2} \iiint\limits_{\text{All Space}} 2E_1(\mathbf{r}) \cdot E_2(\mathbf{r}) d\tau
$$

Interaction Energy:

$$
W^{\mathrm{int}} = \epsilon_0 \iiint\limits_{\mathrm{All\ Space}} E_1(r) \cdot E_2(r) \, d\tau = \frac{q_1 q_2}{16\pi^2 \epsilon_0} \iiint\limits_{\mathrm{All\ Space}} \frac{(r - r_1) \cdot (r - r_2)}{|r - r_1|^3 |r - r_2|^3} \, d\tau
$$

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Example

Find the Interaction energy of two charges, q_1 and q_2 located at r_1 and r_2 , respectively.

Net Electric field at any point r (Superposition Principle)

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Total Configuration Energy:

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$$

$$
= \frac{\epsilon_0}{2} \iiint\limits_{\text{All Space}} |E_1(\mathbf{r})|^2 d\tau + \frac{\epsilon_0}{2} \iiint\limits_{\text{All Space}} |E_2(\mathbf{r})|^2 d\tau + \frac{\epsilon_0}{2} \iiint\limits_{\text{All Space}} 2E_1(\mathbf{r}) \cdot E_2(\mathbf{r}) d\tau
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$$

 \Rightarrow This is exactly the work done by an external agent in bringing q_2 from ∞ to r_2 with q_1 already present at r_1 . **KORKA SERKER ORA**

Interesting Question: Where is the total energy stored?

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 \blacktriangleright ...in the charges?...in the fields?

- \blacktriangleright Interesting Question: Where is the total energy stored?
- \blacktriangleright ...in the charges?...in the fields? No unique answer to that question!
- \blacktriangleright The equation involves integration over all charge distributions:

$$
U_E = \frac{1}{2} \iiint\limits_{V} \rho(\mathbf{r}) V(\mathbf{r}) d^3 r
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 \implies suggests that the energy may be stored in the charges.

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$$

 \implies suggests that the energy may be stored in the fields.

 \blacktriangleright It is conventional to define an ENERGY DENSITY:

$$
u(\mathbf{r}) = \frac{\epsilon_0}{2} \left| \mathbf{E}(\mathbf{r}) \right|^2
$$

 \implies a volume dv will contain Electrostatic Potential Energy equal to $u(r)dv$. **K ロ X K 레 X K 회 X X 회 X 및 X X X X X 전**